ION CYCLOTRON HEATING RATE IN THE SMALL OCTUPOLE

J. C. Sprott

November 2, 1972

PLP 490

Plasma Studies University of Wisconsin

These PLP Reports are informal and preliminary and as such may contain errors not yet eliminated. They are for private circulation only and are not to be further transmitted without consent of the author or major professor.

The spatially averaged cyclotron heating rate in a arbitrary non-uniform magnetic field is given by

$$\frac{d\overline{W}}{dt} = \frac{\pi e \int nE_{\perp}^{2} \delta (B - B_{o}) dV}{2 \int ndV}, \qquad (1)$$

where n is the density, E_{\perp} is the component of the RF electric field perpendicular to the magnetic field, B, and B_0 is the magnetic field at cyclotron resonance. Near the axis of a linear multipole, the B = const. contours are nearly circular and the volume integration can be written

$$dV = r dr d \phi dL$$
.

The heating rate becomes

$$\frac{d\overline{W}}{dt} = \frac{\text{mer}_{o} \int n \ (r_{o}, \phi) \ E_{\perp}^{2} \ (r_{o}, \phi) \ d\phi}{2 \ \frac{dB}{dr} \ r_{o} \iint n \ (r, \phi) \ r \ dr \ d\phi}, \qquad (2)$$

where the subscript o refers to the value of a quantity on the resonance surface. For an octupole, the density along a mod B contour can be approximated by

$$n (r,\phi) \simeq \frac{n_A + n (r,0)}{2} - \frac{n_A - n (r,0)}{2} \cos 4 \phi$$

where the subscript A refers to the value of a quantity on the axis. For a guninjected plasma in the small octupole, the density is known to be approximately

$$n (\psi) \approx n_A^e^{-|\psi - \psi_A|/\psi_1}$$
,

where ψ is the magnetic flux function. Near the axis,

$$B \propto \frac{d\psi}{dr} \propto r^3 \Rightarrow \psi \propto r^4$$

so that

$$n (r,o) = n_A e^{-(r/r_1)^4}$$
.

$$r_o / \frac{dB}{dr} \mid r_o = \frac{r_o^2}{3B_o}$$
.

Finally, it has been shown experimentally that for an electric field produced by a hoop at a distance $r_{\rm H}$ from the axis, the field varies approximately as

$$E_{\perp}^{2}(r) \approx E_{A}^{2} \left[1 + \frac{r^{2}}{r_{H}^{2}} + \frac{2r}{r_{H}} \cos \phi\right]^{-1.6}$$

$$\frac{d\overline{W}}{dt} = \frac{\int_{0}^{\infty} er_{o}^{2} E_{A}^{2} \int_{0}^{\pi} \left[1 + e^{-(r_{o}/r_{1})^{4}} - (1 - e^{-(r_{o}/r_{1})^{4}}) \cos 4\phi\right] \left[1 + \frac{r_{o}^{2}}{r_{H}^{2}} + \frac{2r_{o}}{r_{H}} \cos^{\phi}\right]^{-1.6} d\phi}{\int_{0}^{\infty} \int_{0}^{\pi} \int_{0}^{H} \left[1 + e^{-(r/r_{1})^{4}} - (1 - e^{-(r/r_{1})^{4}}) \cos 4\phi\right] r dr d\phi}.$$

This equation can be written in the form

$$\frac{d\overline{W}}{dt} = \frac{\pi e E_A^2}{6 B_O} F (r_O, r_H, r_1) .$$

For typical conditions (r_H = 6" and r_1 = 4.5"), the function F has been calculated numerically as a function of the resonance zone position (0 < r_0 < r_H), and the result is plotted in the attached figure.

0.1