

REVISED CALCULATION OF OHMIC HEATING RATE
IN THE SMALL OCTUPOLE

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June 1976

PLP 690

Plasma Studies

University of Wisconsin

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I. Introduction

In this paper we present a calculation of the ohmic heating rate for the small octupole with the toroidal magnetic field added. The result may be used in the computer program SIMULT (described in PLP's 505, 556 and 607) to predict the time behavior of spatially averaged plasma parameters. The present calculation is a refinement of the one presented in PLP 643. It is recorded here in full detail for posterity.

II. The Calculation

The ohmic power input is given by the sum over all flux tubes of the average of the power input to each flux tube:

$$dP = \mathcal{V}^2 / \mathcal{R} \quad (1)$$

where \mathcal{V} is the voltage drop along a magnetic field line for one trip the short way around the machine and \mathcal{R} is the electrical resistance along the same path.

Resistance \mathcal{R} is $\int \eta ds / dA$ where η is the local plasma electrical resistivity and ds is the unit length along the path described above and is equal to $B dl / B_p$ where B is the total field strength, B_p is the poloidal field and dl is an element of unit length along a projection of a constant flux surface onto the poloidal plane. The area element dA is the cross-sectional area of the flux tube along which we are calculating voltage drop and resistance. It may be written as $dA = d\Phi_p / B$ where $d\Phi_p$ is the poloidal flux enclosed by neighboring surfaces of a constant poloidal flux. Thus,

$$\mathcal{R} = \langle \eta \rangle \int \frac{B^2}{B_p} \frac{dl}{d\Phi_p}. \quad (2)$$

The voltage drop \mathcal{V} receives contributions from the time-changing toroidal and poloidal fluxes:

$$\mathcal{V} = \dot{\Phi}_T + \frac{\Delta\theta}{2\pi} \dot{\Phi}_p \quad (3)$$

(Cf. Lencioni's thesis PLP 276, eq. 15 & 16, p. 44-45.) The angle $\Delta\theta \equiv \oint \frac{B_T d\ell}{R B_p}$ is the angular advance in the toroidal direction for a field line progressing once around the machine in the poloidal direction.

Substituting equations 2 and 3 into 1 gives:

$$dP = \frac{(\dot{\Phi}_T + \frac{\Delta\theta}{2\pi} \dot{\Phi}_p)^2}{\langle \eta \rangle \oint \frac{B^2}{B_p} \frac{d\ell}{d\Phi_p}} \quad (4)$$

for the ohmic heating in one flux tube. The total power input is thus:

$$P = \frac{\int_{\psi \text{ Limiter}} (\dot{\Phi}_T + \frac{\Delta\theta}{2\pi} \dot{\Phi}_p)^2 d\Phi_p}{\langle \eta \rangle \int_{\psi \text{ Hoop}} \frac{B^2}{B_p} d\ell} \quad (5)$$

The integral in the denominator can be rewritten as:

$$\oint \frac{B^2}{B_p} d\ell = \oint \frac{B_T^2}{B_p} d\ell + \oint B_p d\ell = B_{T0} R_0 \oint \frac{B_T}{R} \frac{d\ell}{B_p} + \oint B_p d\ell =$$

$$B_{T0} R_0 \Delta\theta + \mu_0 I \quad (6)$$

where R_0 is the outer wall major radius, B_{T0} is the value of B_T at R_0 , and I is the current in the hoop or hoops enclosed by the path of integration.

The toroidal flux term is $\dot{\Phi}_T = \int \vec{B}_T \cdot d\vec{A}$. The area element is a shell between two neighboring surfaces of constant poloidal flux and may be written

$$dA = V' \frac{d\psi}{2\pi R_c} \quad (7)$$

where $V' = dV/d\psi$ and V is the volume inside a surface of constant poloidal flux ψ and R_c is the major radius to the center of the machine. Allowing the only approximation in this calculation (except for the averaging of resistivity η), B_T is set equal to its value at the machine's center, B_{TC} . Then

$$\dot{\phi}_T \sim \frac{\dot{B}_{TC}}{2\pi R_C} \int \frac{dV}{d\psi} d\psi = \frac{\dot{B}_{TC}}{2\pi R_C} V \quad (8)$$

The poloidal flux term is $\dot{\phi}_p = \frac{\dot{B}_p}{B_p} \psi_{\text{core}} \frac{(5 + \psi)}{10}$ where the usual small octupole flux-surface labeling convention is followed. This convention is $\psi = -5$ at the hoops and $\psi = +5$ at the walls. Thus also,

$$d\phi_p = \psi_{\text{core}} \frac{d\psi}{10} \quad (9)$$

Designating $\overline{\langle \eta \rangle}$ as the averaged resistivity along a flux tube averaged again over all flux tubes, ohmic heating power becomes

$$P = \frac{1}{\langle \eta \rangle} \int_{\psi_{\text{hoop}}}^{\psi_{\text{limiter}}} \frac{\left[\frac{\dot{B}_{TC}}{2\pi R_C} V - \frac{\Delta\theta}{2\pi} \frac{\dot{B}_p}{B_p} \psi_{\text{core}} \frac{(5 + \psi)^2}{10} \right] \psi_{\text{core}} d\psi}{[B_{TO} R_O \Delta\theta + \mu_O I] 10} \quad (10)$$

The minus sign in the integrand arises from the geometry and sign conventions sketched in Figure 1. The quantities ψ_{core} , B_p and B_T are defined positive in the directions shown. Changing the sign of either B_p or B_T would change the sign of each term of \mathcal{V} leaving the magnitude of \mathcal{V} and P unchanged. Note that the toroidal electric field due to a positive B_p increasing in time is anti-parallel to \vec{B}_T while the poloidal electric field due to a positive B_T increasing in time is in the same direction as \vec{B}_p . The voltage drop along a field line, and thus the ohmic heating, is minimized for \dot{B}_T/B_T and \dot{B}_p/B_p with the same sign.

It is convenient to write the integrand as a function of

$$\alpha = \frac{B_{TO}}{B_{po}} \begin{cases} \text{outer wall} \\ \wedge \text{midplane} \end{cases} \quad (11)$$

Dividing numerator and denominator by B_{TO}^2 , using $B_{TO} R_O = B_{TC} R_C$, and recognizing that I and ψ_{core} are linear with B_{po} ,

$$P = \frac{B_{TO} B_{PO}}{\langle \eta \rangle 10} \frac{\psi_{core}}{B_{PO}} \int_{\psi^{-}hoop}^{\psi limiter} \frac{\left[\frac{\dot{B}_{TO}}{B_{TO}} \frac{R_O}{2\pi R_C^2} V - \frac{\dot{B}_p}{B_p} \frac{\psi_{core}}{B_{po}} \frac{(5 + \psi)}{20\pi} \frac{\Delta\theta}{\alpha} \right]^2 d\psi}{\left[\mu_o \frac{I}{B_{PO}} + R_O \alpha^2 \frac{\Delta\theta}{\alpha} \right]} \quad (12)$$

III. Numerical Details

From PLP 15, Case I, we find

$$\psi_{core} = .3813 B_{po}$$

and $I = 2.995 B_{po}$

for all four hoops.

The function $V'(\psi) = dV/d\psi$ is graphed by Lencioni in his thesis, PLP 276.

An analytic fit (whose form was suggested by Jim Drake) is:

$$V'_p = [-37.45 - 3.75 \psi + 51.7 / \sqrt{\psi_s - \psi}] \times 10^{-3} \text{ m}^3/\text{dory}$$

for the four hoop total in the private flux region and

$$V'_c = [1.336 + 1.782 \psi + 35.52 / \sqrt{\psi - \psi_s}] \times 10^{-3} \text{ m}^3/\text{dory}$$

for the common flux region. The value of ψ at the separatrix is $\psi_s = -1.55$.

Lencioni's graph and our analytic fit are shown in Figure 2. Straightforward integration gives

$$V_p = [51.756 - 37.45 \psi - 1.875 \psi^2 - 103.4 \sqrt{\psi_s - \psi}] \times 10^{-3} \text{ m}^3$$

$$\text{and } V_c = V_p(\psi_s) + (-0.70 + 1.336 \psi + .891 \psi^2 + 71.04 \sqrt{\psi - \psi_s}) \times 10^{-3} \text{ m}^3.$$

These functions are graphed in Figure 3. To check the volume normalization, the above formulae give

$$V_p(\psi_s) = .1053 \text{ m}^3 \text{ and } V_c(\psi_{wall}) = .3160 \text{ m}^3$$

as compared with $V = .3211 \text{ m}^3$ obtained from the machine geometry:

$$V = [L^2 - 2\pi (r_i^2 + r_o^2) - \frac{4}{2} x^2] 2\pi R_C$$

where $L = .356 \text{ m}$ is the length of the tank side, $r_i = 0.22 \text{ m}$ is the inner hoop radius, $r_o = .025 \text{ m}$ is the outer hoop radius and $x^2/2 = (.043)^2/2 \text{ m}^2$ is the area of each corner triangle, and $R_C = .43 \text{ m}$ is the major radius to the center of the tank.

The angular transform $\Delta\theta$ in degrees is graphed as a function of ψ by Poukey and Howard in PLP 91 and tabulated by Howard in PLP 106. An analytic fit in the private flux region is

$$\Delta\theta_p/\alpha = -3.101 + 2.41 \psi + 43/\sqrt{\psi_S - \psi}$$

and in the common flux region

$$\Delta\theta_c/\alpha = 5 \quad 127.3/\sqrt{\psi - \psi_S}$$

Both graphs and analytic fits are shown in Figure 4.

The square in the integrand of Eq. 12 may be expanded so that the heating power may be written:

$$P = \frac{\alpha B_{TO} B_{PO}}{2\pi} \frac{\psi_{core}}{10 B_{PO}} \left[A \left(\frac{B_{TO}}{B_{TO}} \right)^2 I_A + C \left(\frac{B_P B_{TO}}{B_P B_{TO}} \right) I_C + D \left(\frac{B_P}{B_P} \right)^2 I_D \right] \quad (13)$$

where

$$A = \left(\frac{R_O}{R_C} \right)^2 \frac{1}{(2\pi R_C)^2} = \left(\frac{.61}{.43 \times 2\pi \times .43} \right)^2 = .275$$

$$C = \left(\frac{R_O}{R_C} \right) \frac{1}{2\pi R_C} \frac{\psi_{core}}{10 B_{PO}} = \frac{.61 \times .3813}{.43 \times 2\pi^2 \times .43 \times 10} = 6.373 \times 10^{-3}$$

$$D = \left(\frac{\psi_{core}}{B_P} \frac{1}{20\pi} \right)^2 = \left(\frac{.3813}{20\pi} \right)^2 = 3.683 \times 10^{-5}$$

and

$$I_A = 4 \int_{\psi_H}^{\psi_S} \frac{(V_P/4)^2 d\psi}{D_P} + \int_{\psi_S}^{\psi_L} \frac{V_C^2}{D_C} d\psi$$

$$I_C = 4 \int_{\psi_H}^{\psi_S} \frac{d\psi (\Delta\theta)}{D_P} \int_{\psi_S}^{\psi_L} \frac{V_C (\Delta\theta_c/\alpha) (5)}{D_C}$$

$$I_D = 4 \int_{\psi_H}^{\psi_S} \frac{(\Delta\theta_p/\alpha)^2 [(5 + \Psi)/4]^2 d\psi}{D_P} + \int_{\psi_S}^{\psi_L} \frac{(\Delta\theta_c/\alpha)^2 (5 + \psi)^2 d\psi}{D_C}$$

where Ψ_H , Ψ_S and Ψ_L designate the locations in Ψ -space of the hoop surface, the separatrix and the limiter, and

$$D_P = \mu_0 I/4 + R_0 \alpha^2 \Delta\theta_p/\alpha = .7487 + .61 \alpha^2 \Delta\theta_p/\alpha$$

$$D_C = \mu_0 I + R_0 \alpha^2 \Delta\theta_c/\alpha = 2.995 + .61 \alpha^2 \Delta\theta_c/\alpha.$$

The integrals I_A , I_C and I_D are done numerically. Functions of the form $K_1/(K_2 + \alpha^2)$ fit these integrals extremely well. Combining the numerical coefficients of the integrals with the numerical coefficients of Equation 13 gives the final result for ohmic heating power density.

$$\frac{P}{Vol.} = \frac{\alpha B_{TO} B_{PO} 10^{-3}}{\langle \eta \rangle} \left[\frac{C_1}{(C_2 + \alpha^2)} \left(\frac{\dot{B}_{TO}}{B_{TO}} \right)^2 - \frac{C_3}{(C_4 + \alpha^2)} \left(\frac{\dot{B}_{TO} \dot{B}_{PO}}{B_{TO} B_{PO}} \right) + \frac{C_5}{(C_6 + \alpha^2)} \left(\frac{\dot{B}_{PO}}{B_{PO}} \right)^2 \right] \quad (14)$$

all in MKS units where the volume is taken to be $V(\psi_L)$. The constants C_i , $i = 1 - 6$, are tabulated in Table I for four values of ψ_L . The upper table is for α defined as B_{TO}/B_{PO} . The lower table is for α defined as B_{PO}/B_{TO} . The same formula for power density (Eq. 14) holds for either definition of α provided the correct coefficients C_i are used in each case.

This result differs slightly from the result of PLP 643 in both magnitude and α dependence. For typical conditions ($\alpha \approx 1$), the present result predicts a slightly smaller ohmic heating.

IV. Analysis of the Result

It is now interesting to calculate the ohmic heating power as a function of time for the existing experiment (see PLP's 623 and 674 for description of experiment).

The heating rate is a function of time because B_{PO} , B_{TO} and α are different functions of time.

The plasma quantities such as density n_e , electron temperature T_e and degree of ionization affect the ohmic heating rate because they affect the resistivity. The resistivity on axis has been found experimentally to be the sum

$$\eta = \eta_{sp} + \eta_L + \eta_n$$

where η_{sp} is the classical resistivity for $z = 1$ calculated by Spitzer:

$$\eta_{sp} = 5.21 \times 10^{-5} z \ln \Lambda / T_e^{3/2} \quad \Omega - m$$

where $\ln \Lambda \approx 13$ is the Coulomb logarithm for our usual plasma parameters and η_L is the resistivity measured at low densities by both Lencioni (see his thesis, PLP 276) and us (PLP 674):

$$\eta_L = 2.5 \times 10^7 \sqrt{T_e} / n_e \quad \Omega - m$$

and η_n is the resistivity due to electron collisions with neutral hydrogen:

$$\eta_n = .6 \times 10^{-5} n_H / n_e \quad \Omega - m.$$

For a nearly ionized plasma ($n_H \ll n_e$), η_n may be neglected.

Ohmic heating power as a function of time was calculated from Eq. 14 for moderate B_P and B_T capacitor bank voltages and the easily attainable density $n_e = 8 \times 10^{11} / \text{cm}^3$ for several values of electron temperature T_e . For simplicity, n_e and T_e were assumed constant in time. The results are graphed in Figure 5. The highest heating rate is attained for that value of T_e which minimizes resistivity for the given value of n_e . As the graph in Figure 7a shows, η is a minimum for $n_e \approx 10^{12} / \text{cm}^3$ when $T_e = 10$ eV. A higher plasma density would allow η to be minimized and hence heating power to be maximized at a higher value of T_e .

Note that the heating rate drops by roughly a factor of 100 during the experiment. This decline is due to the relative time scales and magnitudes of B_{TO} and B_{PO} . In Figures 6 - 11 the time behavior of ohmic heating power is explored through the existing available ranges of B_P and B_T for the typical plasma parameters

$n_e = 8 \times 10^{11}/\text{cm}^3$ and $T_e = 10$ eV. Because B_T has a period about three times shorter than B_p in our usual experiment, the \dot{B}_T^2 term dominates heating during the B_T rise time. Thus, the initial heating rate varies approximately as B_T^2 , as shown in Figure 6. But changing B_T makes scant difference during the latter part of the experiment where heating is due to the small \dot{B}_p^2 term. Figure 7 shows the contribution from each term in Eq. 14 for a large B_T and moderate B_p . The heating rate begins high but drops very fast. After B_T peaks, the $\dot{B}_p \dot{B}_T$ cross term becomes negative and nearly cancels the contributions of the positive terms near the end of the experiment.

In Figure 8 the effect of varying B_p is displayed. Since the \dot{B}_T^2 term contributes most of the early power, the initial rate is hardly affected by changing B_p . But increasing B_p allows power to fall less sharply with time. Time behavior of the three terms for high B_p and moderate B_T is graphed in Figure 9.

As shown in Figure 10, the highest heating rates result from the largest possible B_p and B_T both. For equal capacitor voltages on each bank, power increases as the square of the fields. Figure 11 shows that in this case also the \dot{B}_T^2 term is initially very large. The \dot{B}_p^2 term is smaller than the initial heating rate, but dominates when B_T is peaking. Near the end of the experiment, the negative cross term unfortunately nearly cancels the positive contributions.

With the ability to calculate ohmic heating rates in the small octupole, we may investigate possible modifications of the experiment to try to increase magnitude and/or duration of ohmic power input to the plasma. The obvious variables besides magnitudes discussed above are relative timing and pulse lengths. A full study is left for the future, but Figures 12 and 13 show the effect of crowbarring B_T at its peak to give an e-folding decay time of 5 ms. Figure 12 compares heating power with and without crowbarring. After the fast drop of the initial high rate, the power is held between 55 and 100 kW for almost 2 ms by crowbarring B_T . The

increased power is achieved by making the $\dot{B}_T \dot{B}_P$ cross term negligible after the B_T peak so that the full effect of the \dot{B}_P^2 term can be realized. This is shown in Figure 13.

Other ways to increase the ohmic heating input power and duration no doubt exist within the capabilities of the present experiment or with moderate modifications. Using the analytical form of Eq. 14 for input power will enable us to investigate various possibilities. Their effect on the experiment, however, also depends on the heat loss terms which are not yet fully understood.

V. Using the Result in SIMULT

The formula for ohmic heating input power (Eq. 14 with $\psi_L = 1$) was included in subroutine OHMIC which program SIMULT calls whenever there is a toroidal field. The program was run with the same parameters as those in PLP 643: $B_{T \text{ MAX}}$ on axis = 1.25 kG (2.0 kV on B_T bank), $B_{PO} = 0.88$ kG (2.0 kV on B_P bank), $p = 1 \times 10^{-5}$ torr (gauge), $Z_{\text{eff}} = 1$ with 1 kW of 2.45 GHz microwaves turned off and B_T turned on at 3.5 msec into the 5 msec poloidal field pulse. Since PLP 643, SIMULT has been refined in a number of other areas, the most significant of which are:

- 1) A more realistic time and density dependent microwave heating rate has been added.
- 2) The neutral reflux term has been modified to include ion impact desorption.
- 3) Losses due to the finite ion gyroradius have been included.
- 4) The classical diffusion term has been refined.

The results along with a current listing of SIMULT are included in the appendix. Note that the maximum density is $8.1 \times 10^{11} \text{ cm}^{-3}$ (vs $7.4 \times 10^{11} \text{ cm}^{-3}$ in PLP 643), the maximum electron temperature is 75 eV (vs 100 eV in PLP 643). The density is in reasonable agreement with the experiment (note that the plasma is 96% burned out), but the temperature is still higher than what is measured in the experiment (10 - 20 eV). The discrepancy may be due to a number of factors that we intend to investigate:

1) The resistivity measured on axis may not be typical of field lines well off the separatrix.

2) Impurity radiation may carry out much of the electron energy.

3) There may be additional particle or energy losses resulting from instabilities or waves, in the presence of the ohmic heating.

Table I. Coefficients for the Ohmic Heating Power Formula for the Small Octupole.

$$\alpha \equiv [B_{T0}/B_{pol}|_{\text{outer wall}} \wedge \text{midplane}]$$

Ψ_L	c_1	c_2	c_3	c_4	c_5	c_6
1.0	2.40	1.69	4.13	1.51	1.99	1.04
2.0	3.89	1.88	6.18	1.72	2.69	1.26
3.0	5.49	1.99	8.58	1.86	3.56	1.44
5.0	8.51	2.01	13.9	1.94	5.85	1.68

$$\alpha \equiv [B_{p0}/B_{T0}|_{\text{outer wall}} \wedge \text{midplane}]$$

Ψ_L	c_1	c_2	c_3	c_4	c_5	c_6
1.0	1.42	.591	2.74	.663	1.92	.962
2.0	2.06	.533	3.60	.583	2.14	.794
3.0	2.76	.503	4.62	.539	2.47	.693
5.0	4.24	.498	7.18	.516	3.48	.595

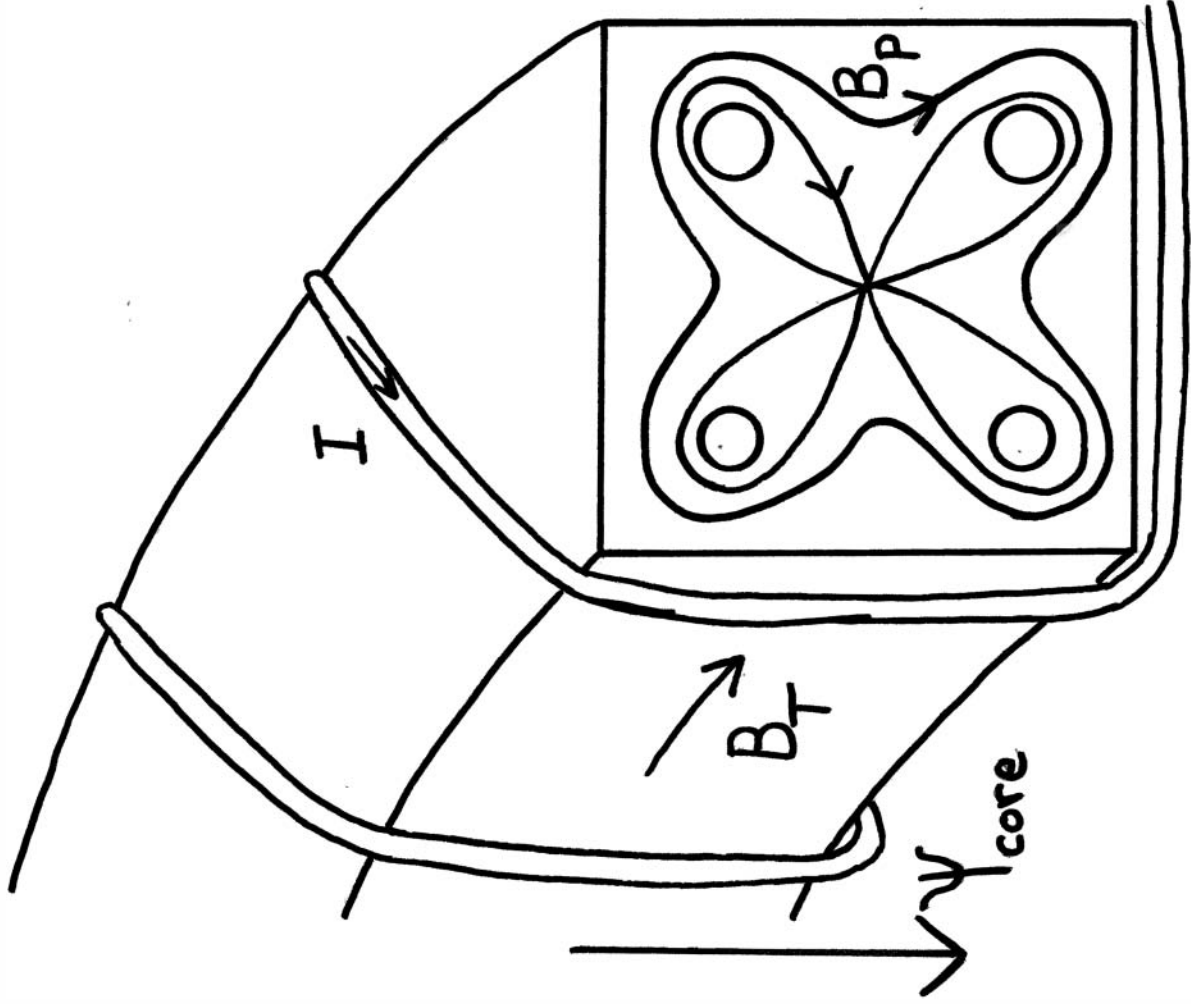
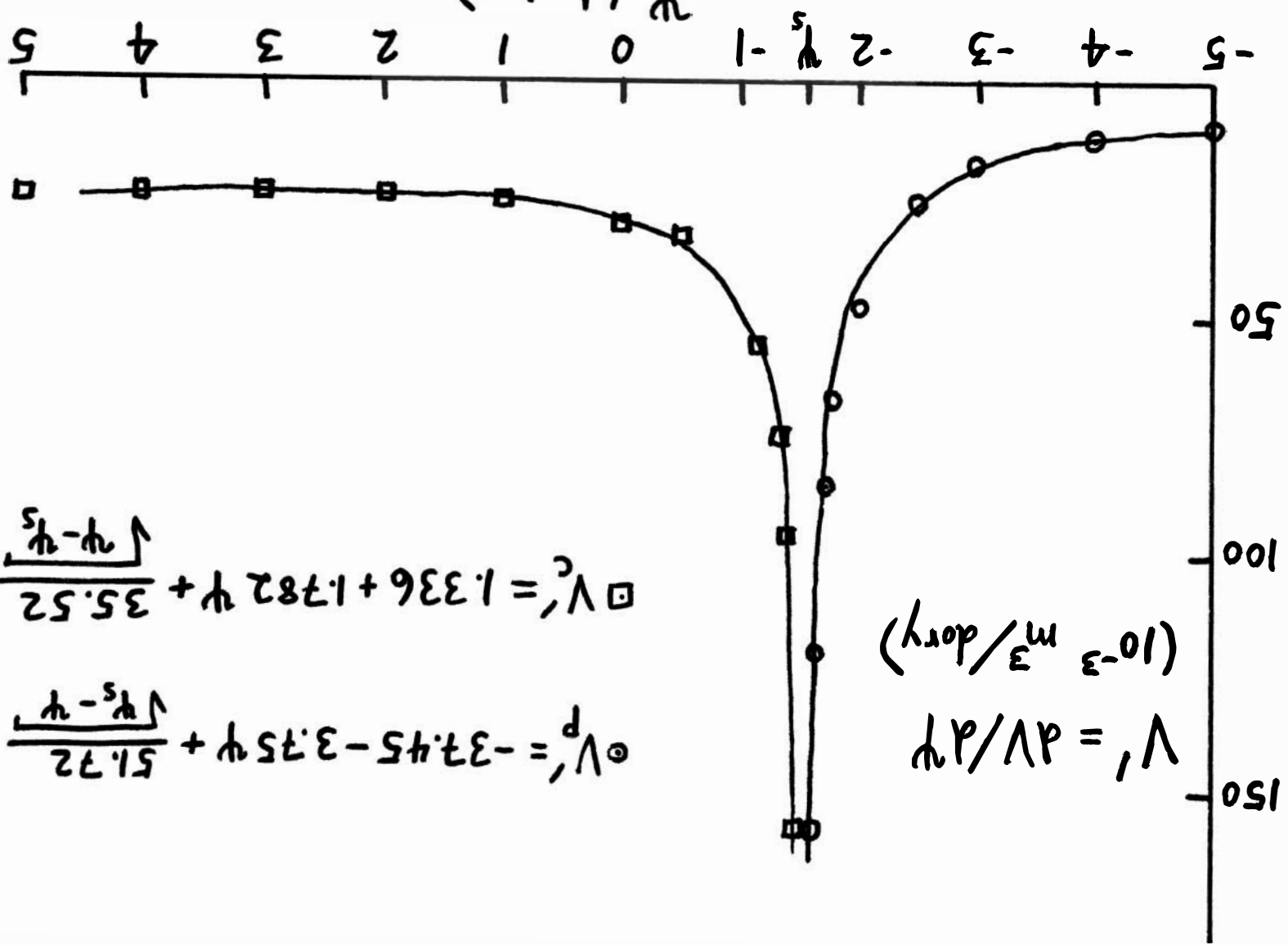


Figure 1. Magnetic Field Geometry and Sign Conventions.

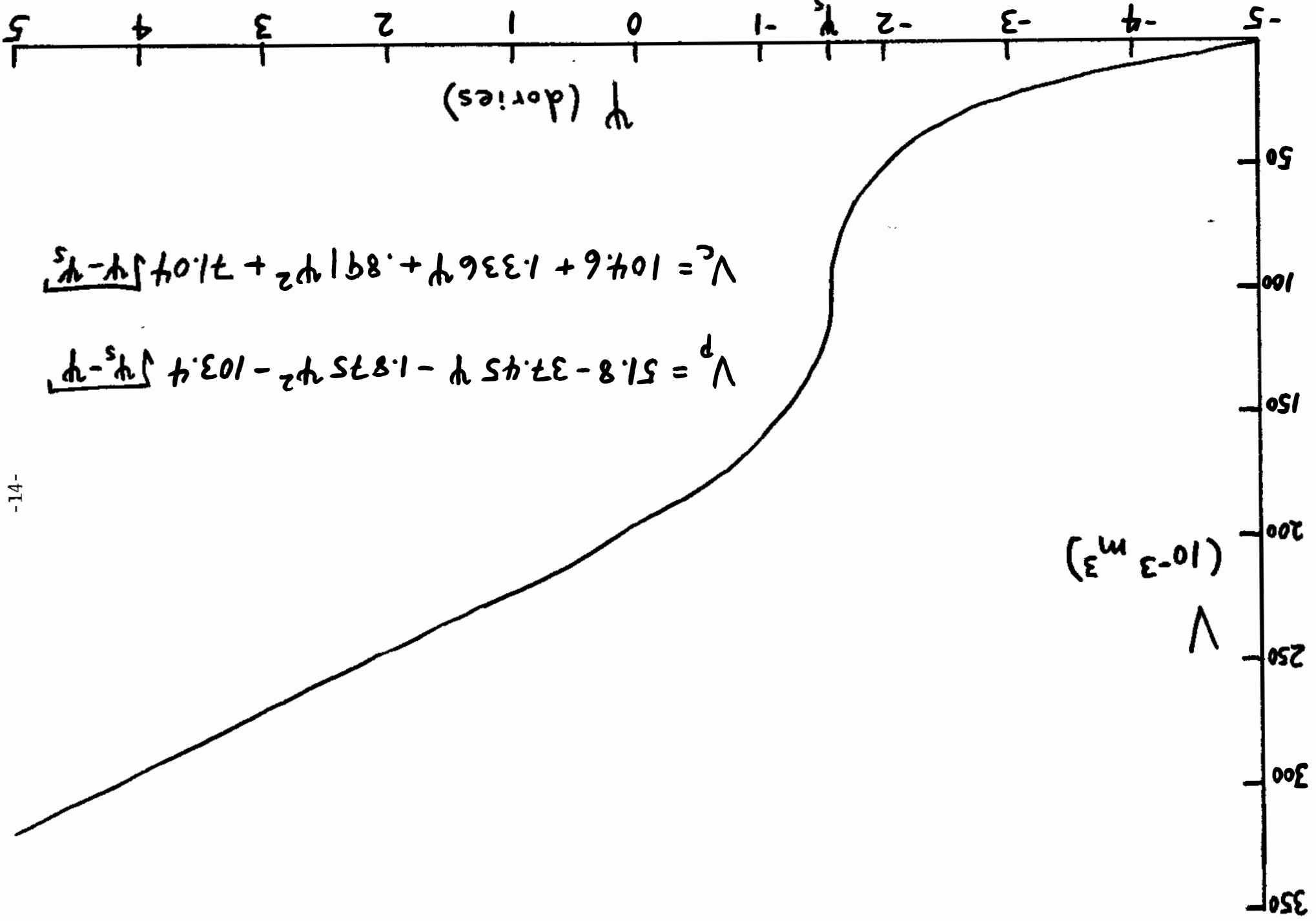
Figure 2. $V'(\psi)$



$$\square V_p^2 = 1.336 + 1.782\psi + \frac{\sqrt{\psi - \psi_s}}{35.52}$$

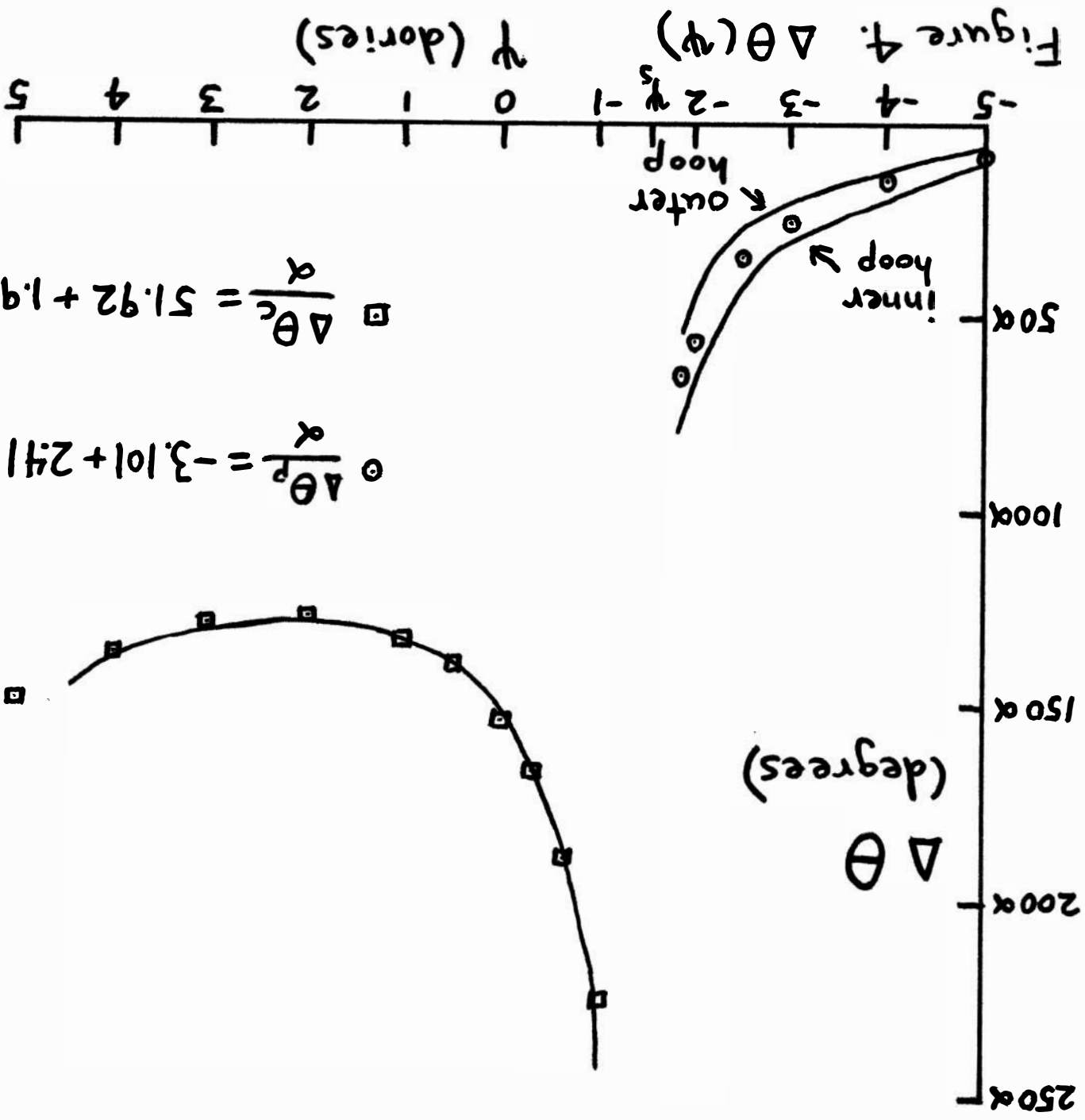
$$\circ V_p^1 = -37.45 - 3.75\psi + \frac{\sqrt{\psi_s - \psi}}{51.72}$$

Figure 3. $V(\psi)$



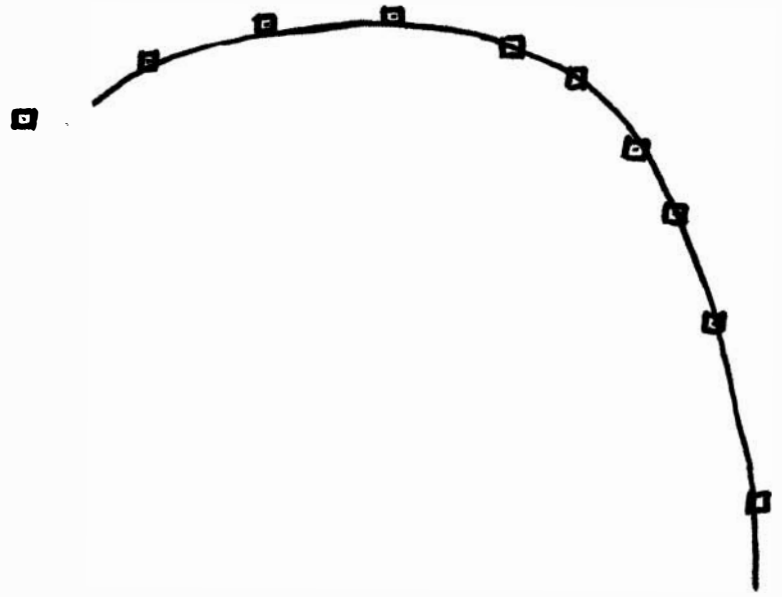
$$V_c = 104.6 + 1.336\psi + 1.891\psi^2 + 71.04\sqrt{\psi - \psi^5}$$

$$V_p = 51.8 - 37.45\psi - 1.875\psi^2 - 103.4\sqrt{\psi - \psi^5}$$



$$\Delta\theta_P = -3.101 + 2.41\psi + \frac{\sqrt{\psi - \psi_s}}{4.3}$$

$$\Delta\theta_C = 51.92 + 1.909\psi^2 + \frac{\sqrt{\psi - \psi_s}}{127.3}$$



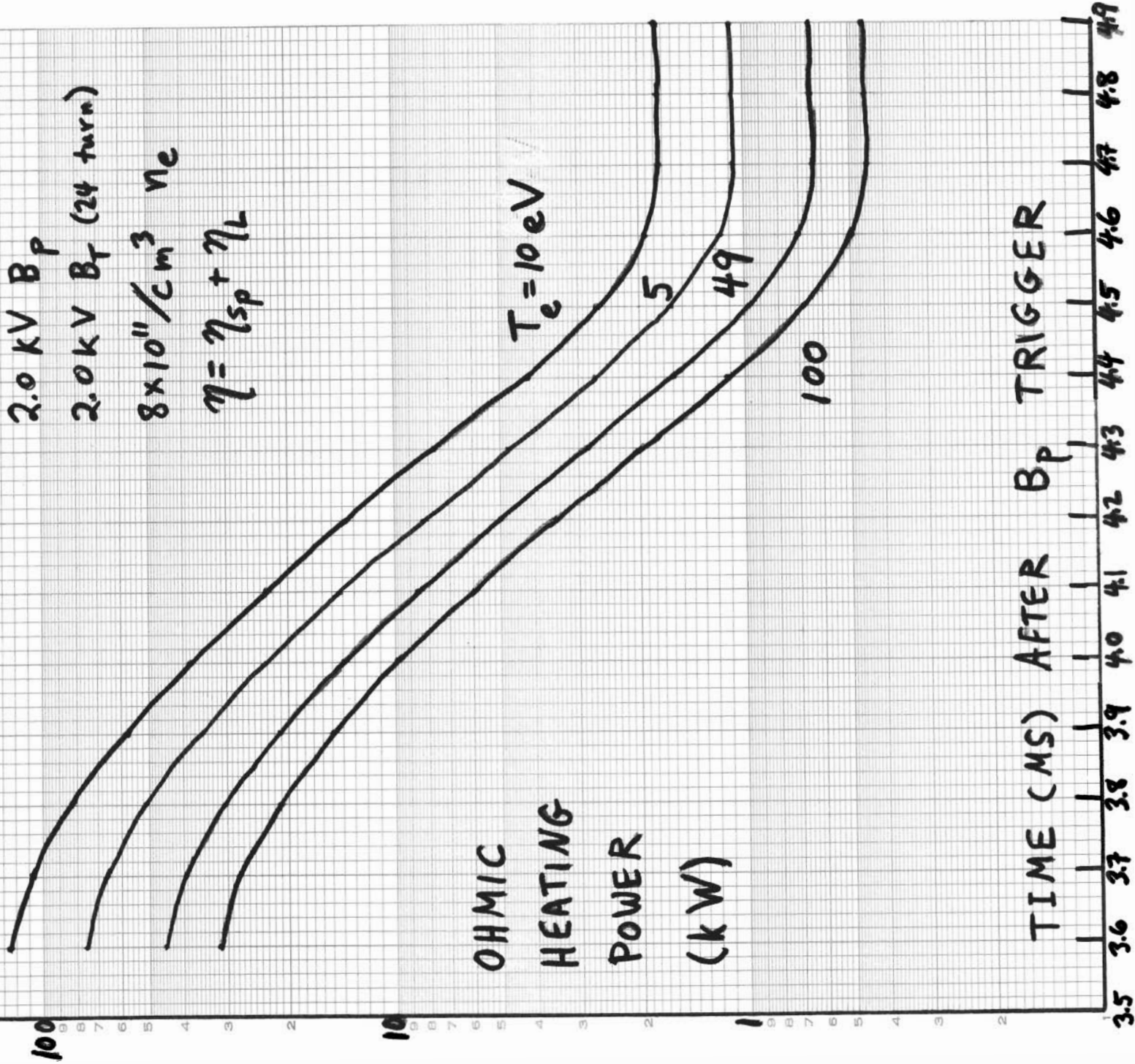


Figure 5. OHMIC HEATING POWER VERSUS TIME FOR VARIOUS T_e CORRESPONDING TO VARIOUS $\langle \eta \rangle$.

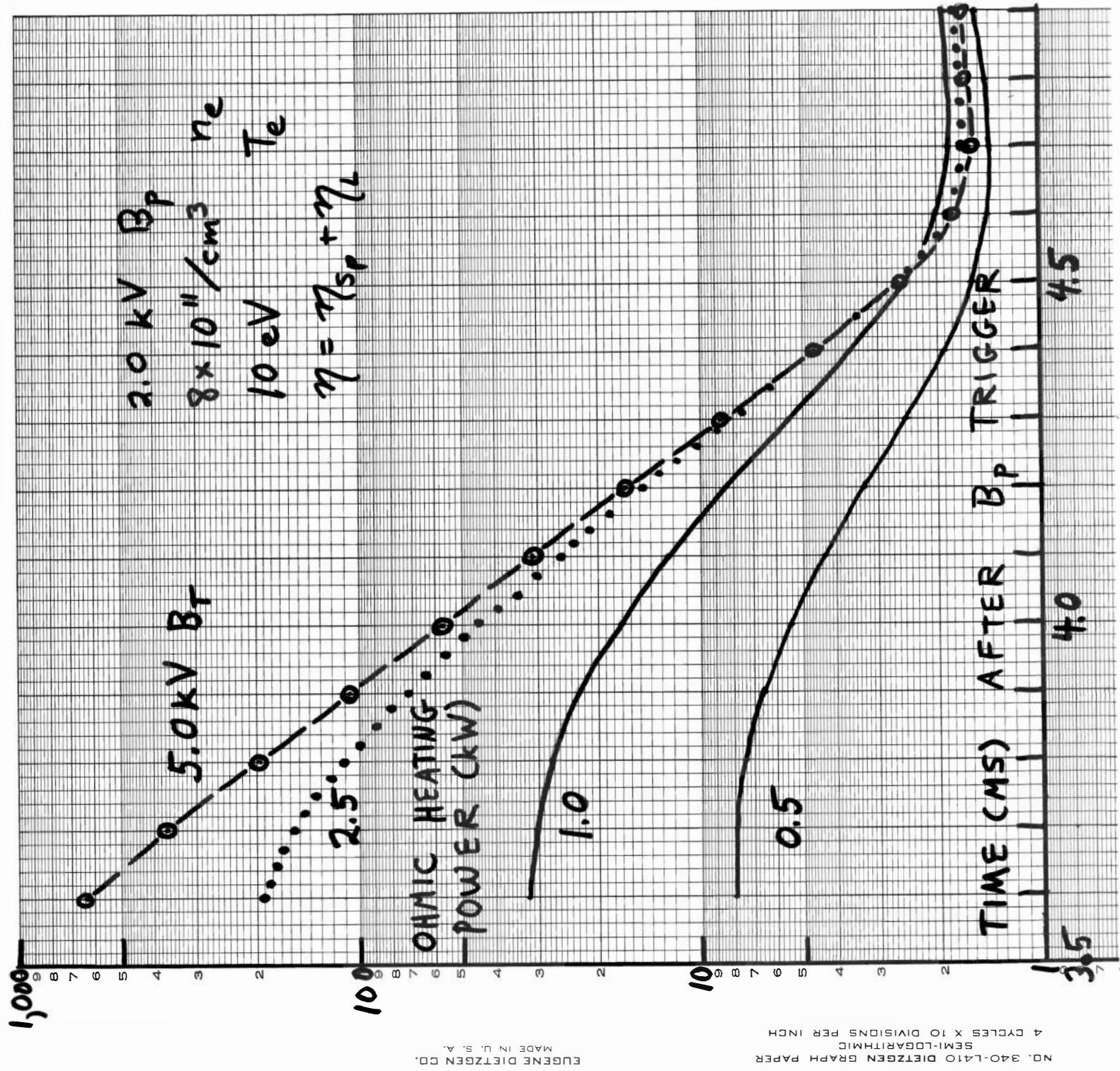


Figure 6. OHMIC HEATING POWER VERSUS TIME FOR VARIOUS B_T

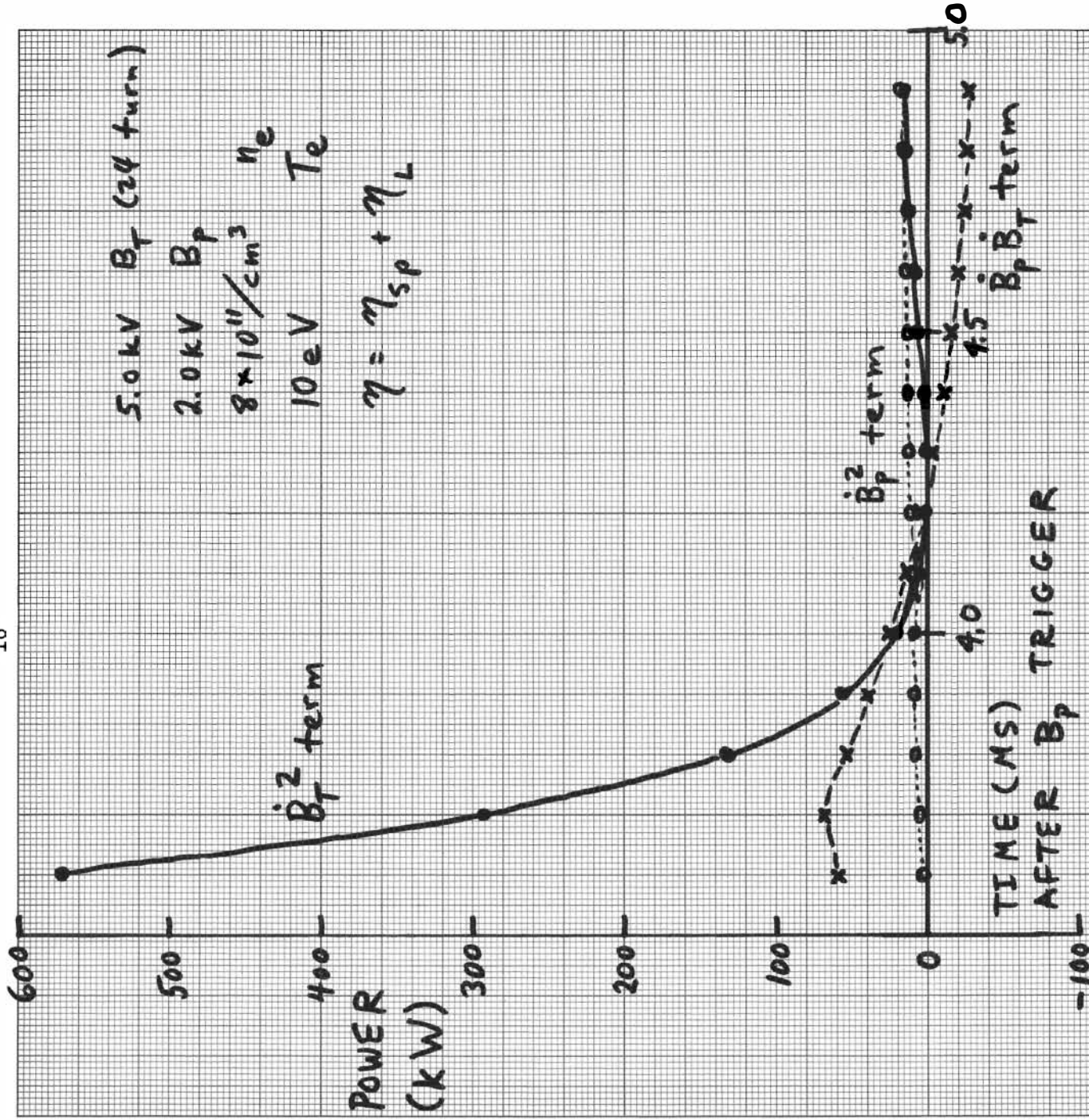
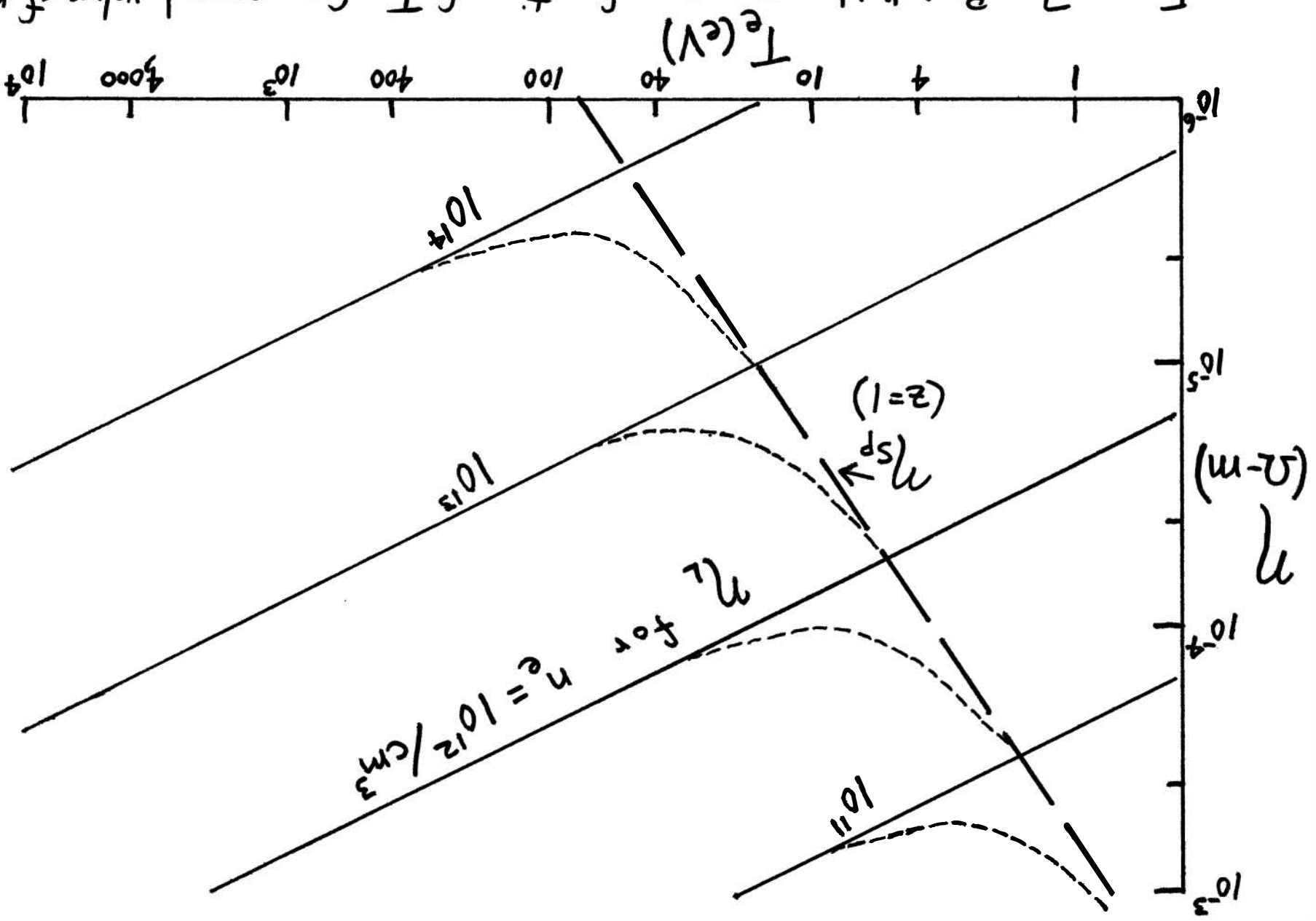


Figure 7. THE THREE CONTRIBUTIONS TO OHMIC HEATING POWER VERSUS TIME FOR LARGE B_T AND MODERATE B_p .

Figure 7a. Resistivity η as a function of T_e for several values of n_e .



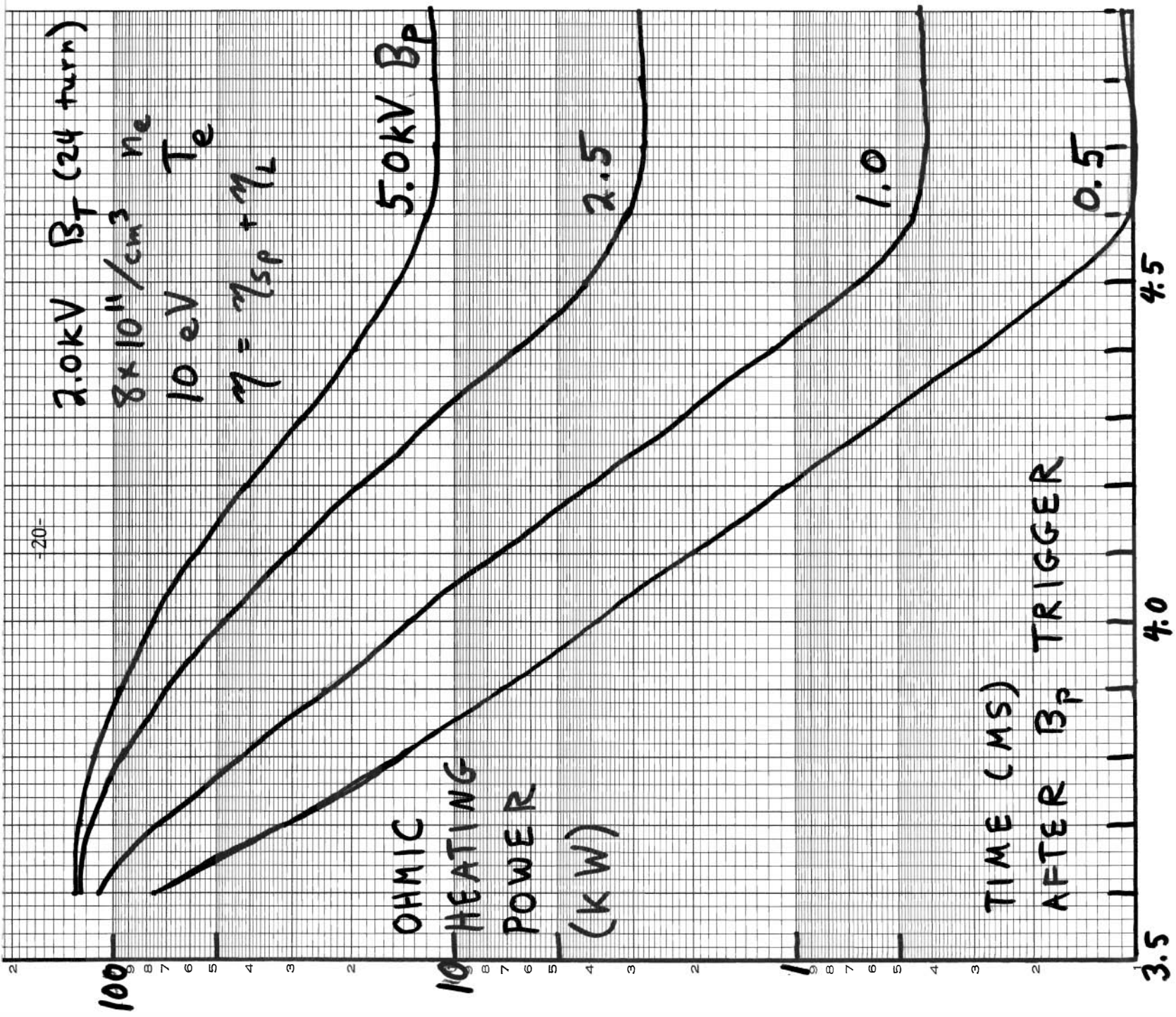
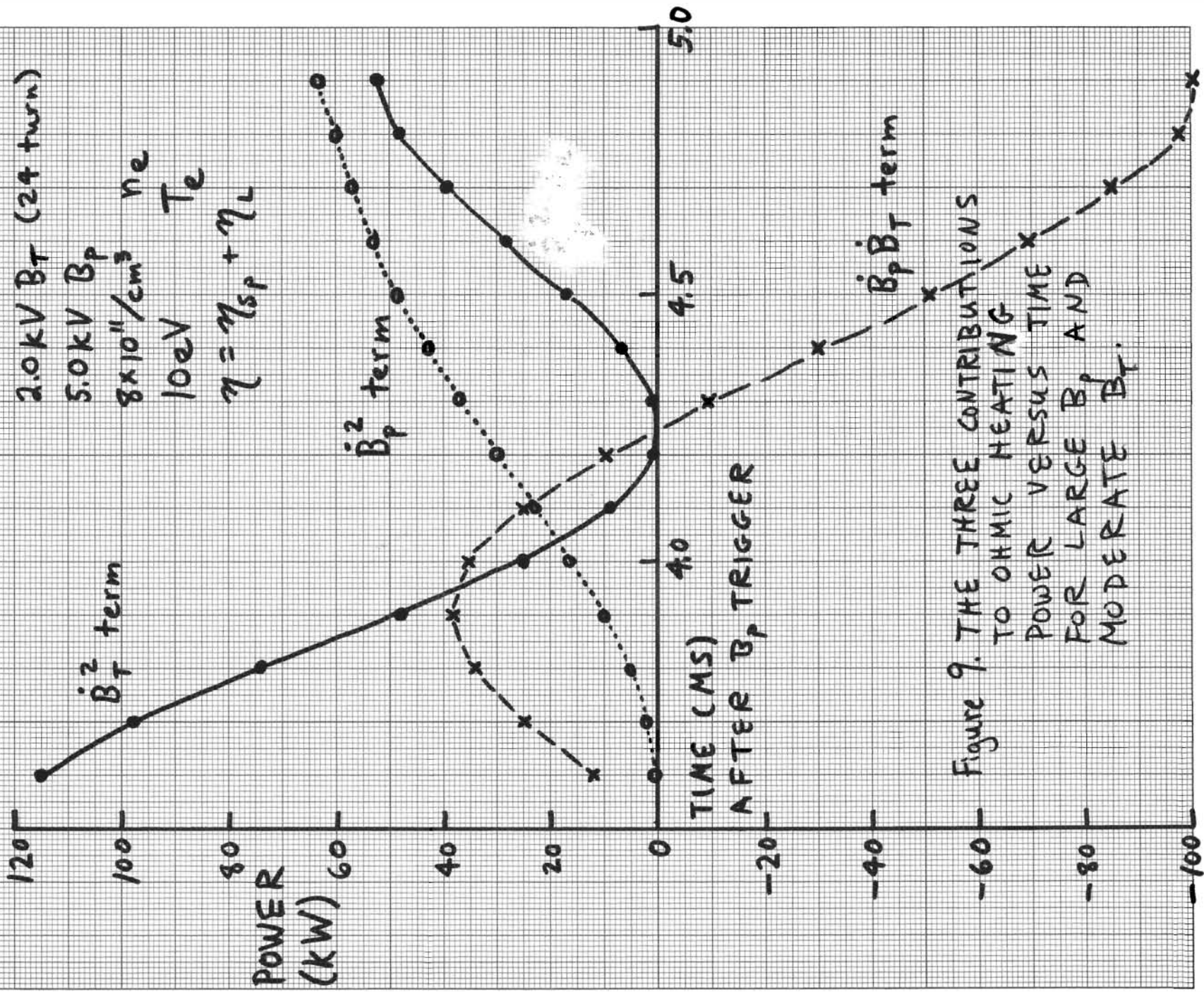


Figure 8. OHMIC HEATING POWER VERSUS TIME FOR VARIOUS B_P



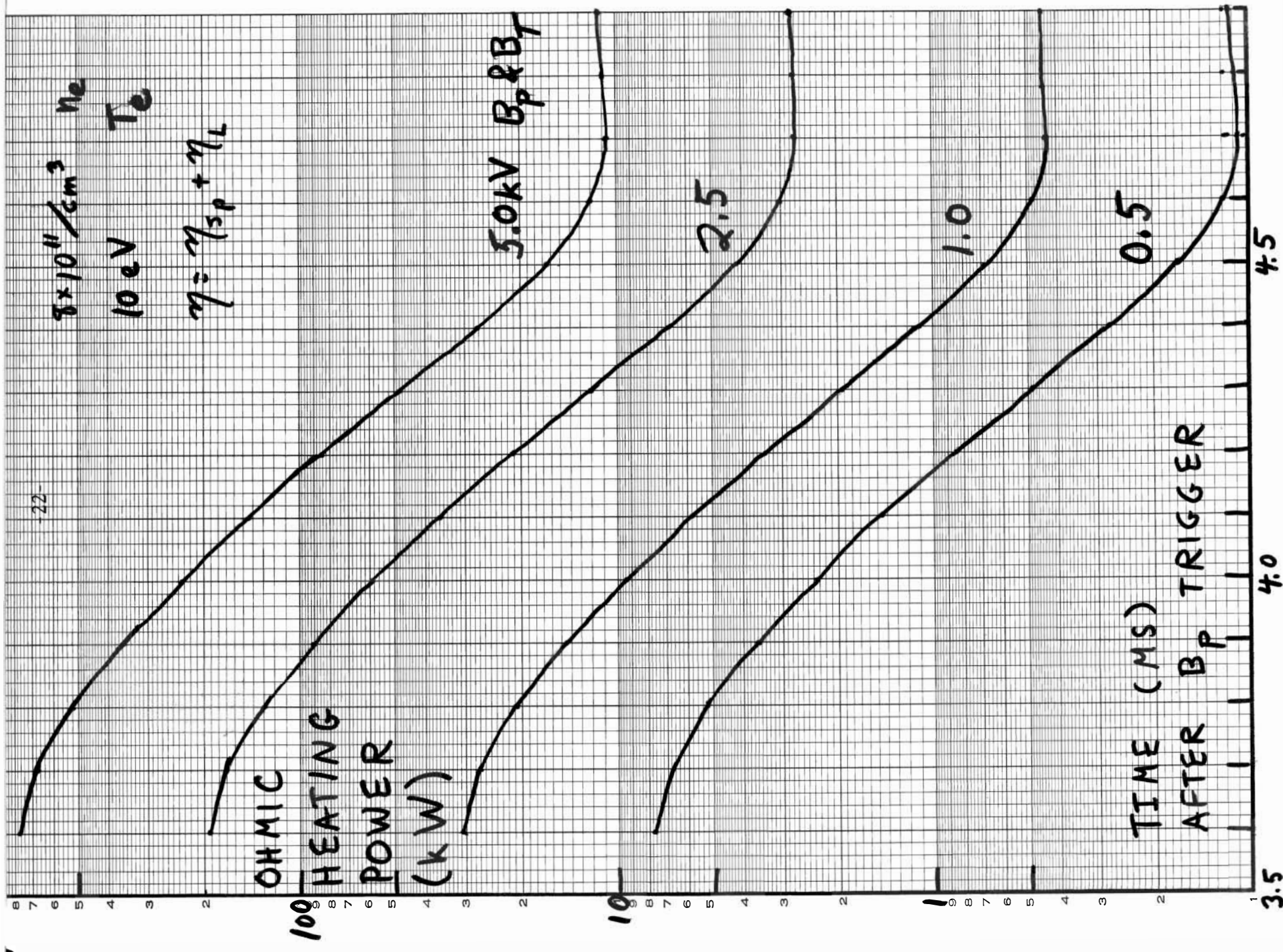


Figure 10. OHMIC HEATING POWER VERSUS TIME FOR VARIOUS B_p AND B_T BUT SAME RATIO OF PEAK VALUES

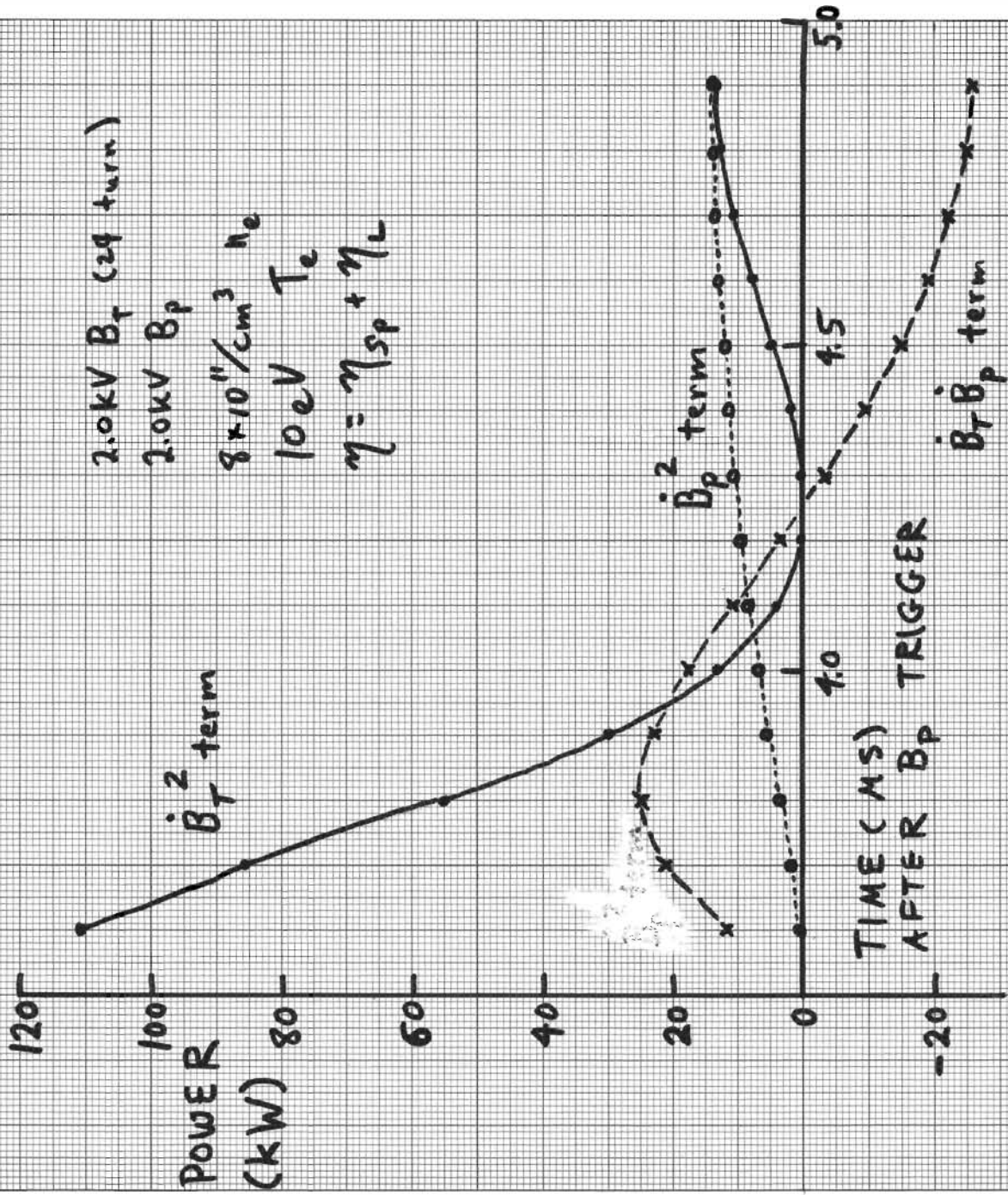


Figure 11. THE THREE CONTRIBUTIONS TO OHMIC HEATING POWER VERSUS TIME FOR SAME CHARGE OF B_T AND B_p CAPACITOR BANKS,

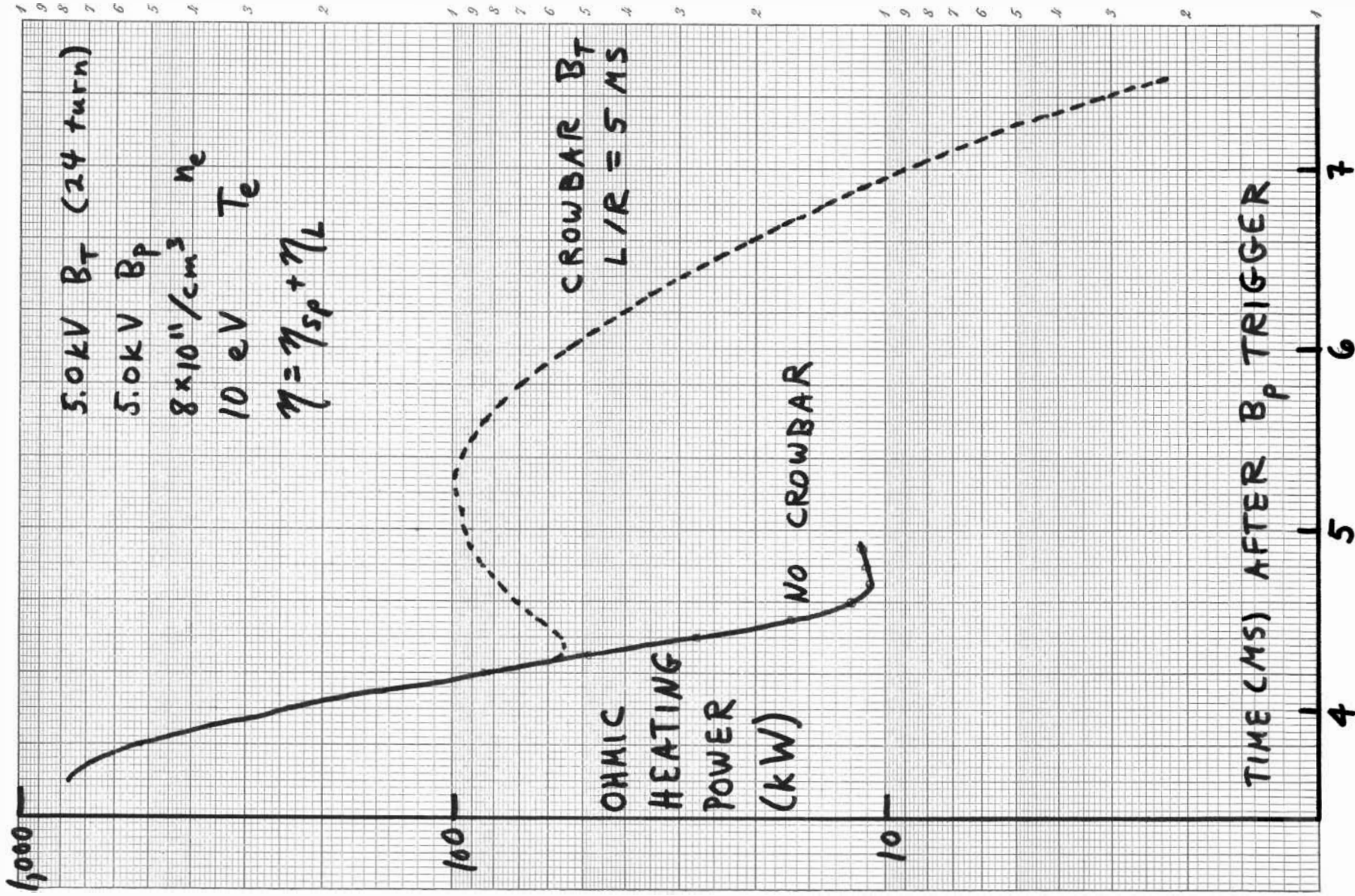


Figure 12. OHMIC HEATING POWER VERSUS TIME FOR CROWBARRED AND UN-CROWBARRED B_T .



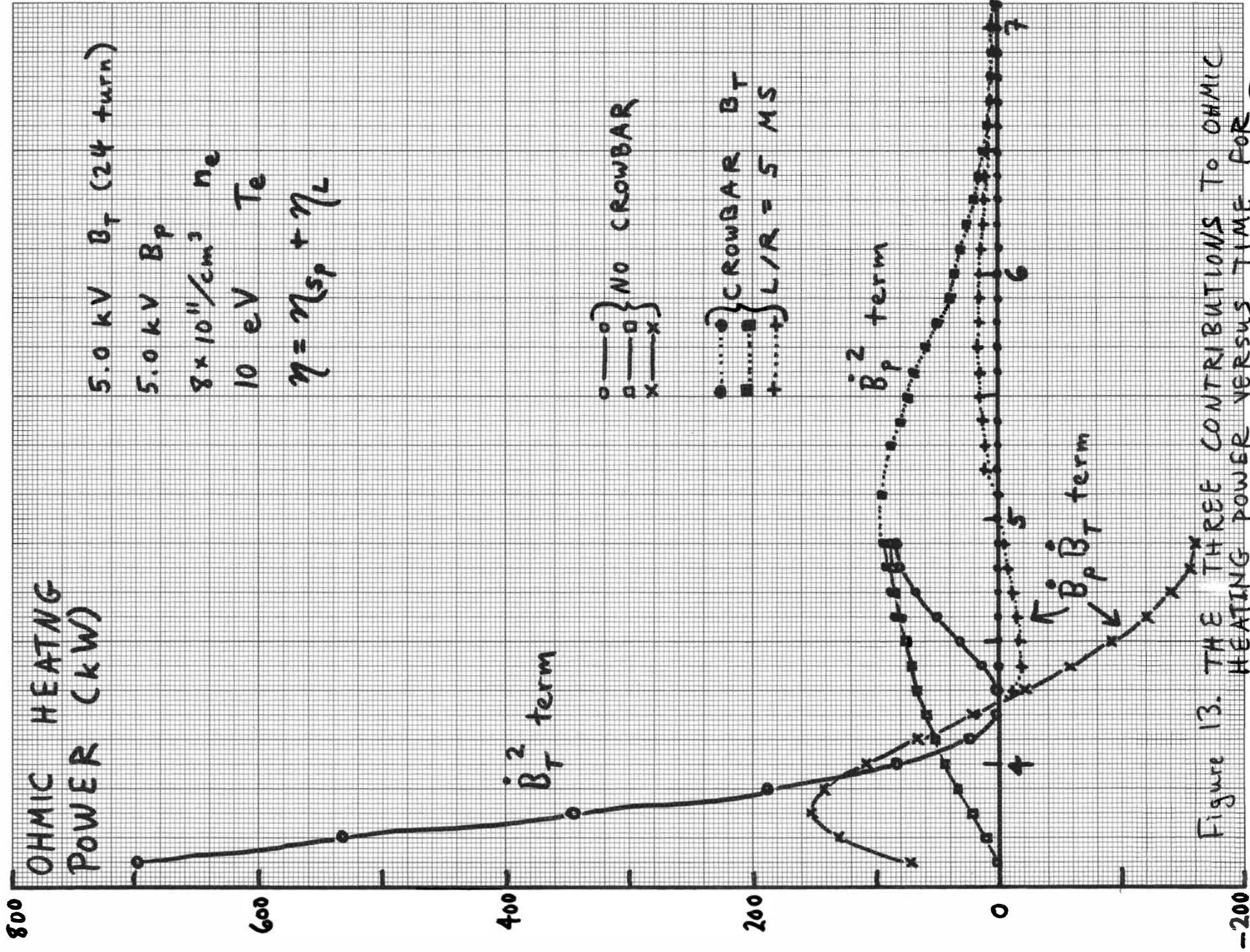


Figure 13. THE THREE CONTRIBUTIONS TO OHMIC HEATING POWER VERSUS TIME FOR CROWBARRED AND UNCROWBARRED B_T

APPENDIX: SIMULT PROGRAM AND RESULTS

UNIVAC 1110 TIME SHARING EXEC. --- MULTI-PROCESSOR SYSTEM --- VER. 31.6
 PROJECT: 102980 USER: 4126810219 FILE: PR000C00295 MACC: 31.6 PART: 00

DTT, 2980, 4126810219, 2M

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SIMULT
ACC 1.178-06/08/76-14143119 SIMULT
1 C PROGRAM SIMULT J.C. SPROTT - JAN 25, 1973
2 C EXTERNAL OHMIC, HYCRH, HECRH
3 C DIMENSION T(101), DE(101), TE(101), TI(101)
4 C DIMENSION DN(101), BP(101), PP(101), SJ(101)
5 C DIMENSION SD(5), STE(6), STI(4), D(5), PE(6), PI(4)
6 C DEFINE FUNCTIONS - OCTUPOLE
7 C D1 IS DN/DT DUE TO IONIZATION
8 C D1(DENS, DNEUT, TE) = 371.0 * DENS * DNEUT * SQRT(TE) * EXP(-15.6/TE) * (TE / (20.
9 C 20 * TE + 15.6) + LOG(1.5625 + 0.1 * TE)) / (TE + 15.6)
10 C D2 IS DN/DT DUE TO DIFFUSION
11 C D2(DENS, TE) = DENS * (DENS / SQRT(TE) + 0.026 * DNEUT * TE) / B / B / A / A
12 C D3 IS DN/DT DUE TO OBSTACLE LOSSES
13 C D3(DENS, TE) = 2.0E5 * DENS * AO * SQRT(TE + TIA) / A / A / AL
14 C D4 IS DN/DT DUE TO FIELD DECAY
15 C D4(DENS) = 0.5 * DENS * AMAX(0.0, BOLD - B) / BOLD / DT
16 C D5 IS DN/DT DUE TO FINITE ION GYORADIUS
17 C D5(DENS, TI) = 1300.0 * DENS * 2 / TI * 1.5 * EXP(AMINI(2.4 * A * A * B * B / TI, 88.0))
18 C PE1 IS DUE/DT DUE TO MICROWAVES
19 C PE1(P) = 2.0E9 * P * G * DEA / (G * DEA + 12.35 * P * G) / A / A / AL
20 C PE2 IS DUE/DT DUE TO ION COLLISIONS
21 C PE2(DENS, TE, TI) = 2.3 * DENS * 2 * (TE - TI) * LOG(5.2E11 * TE * 3 / ABS(DENS)
22 C 2 / (40.0 + TE)) / TE * 1.5
23 C PE3 IS DUE/DT DUE TO EXCITATION
24 C PE3(DENS, DNEUT, TE) = 29.1 * D1(DENS, DNEUT, TE) * EXP(6.96 / (TE + 0.1))
25 C PE4 IS DUE/DT DUE TO BREMSSTRAHLUNG
26 C PE4(DENS, TE) = 1.0E-4 * DENS * DENS * SQRT(TE)
27 C PE5 IS DUE/DT DUE TO SYNCHROTRON RADIATION
28 C PE5(DENS, TE) = 3.87E-3 * DENS * B * B * TE * (1.0 + TE / 2.04E5)
29 C PE6 IS DUE/DT DUE TO THERMAL CONDUCTION
30 C PE6(DENS, TE) = (2.5 * D2(DENS, TE) + 2.0 * D3(DENS, TE) + 1.5 * D4(DENS) + 1.5 * D5(
31 C 2DENS, TIA)) * (TE - TWALL)
32 C PI3 IS DUE/DT DUE TO CHARGE EXCHANGE
33 C PI3(DENS, DNEUT, TI) = 0.732 * DENS * DNEUT * SQRT(TI) * (TI - TWALL)
34 C 2 * (1.0 + 0.00585 * TI * 1.5) / EXP(0.0582 * SQRT(TI))
35 C PI6 IS DUE/DT DUE TO THERMAL CONDUCTION
36 C PI6(DENS, TE, TI) = (2.5 * D2(DENS, TE) + 2.0 * D3(DENS, TE) + 1.5 * D4(DENS)) * (TI
37 C 2 - TWALL) + (2.4 * A * A * B * B - 1.5 * TWALL) * D5(DENS, TI)
38 C PI7 IS DUE/DT DUE TO IONS BORN AT TEMP OF WALL
39 C PI7(DENS, DNEUT, TE) = 1.5 * D1(DENS, DNEUT, TE) * TWALL
40 C PI8 IS DUE/DT DUE TO ELASTIC COLLISIONS WITH NEUTRALS
41 C PI8(DENS, DNEUT, TI) = 10.0 * DENS * DNEUT * (TI - TWALL) / (570. + TI * 2.4) * 29
42 C SPECIFY PARAMETERS
43 C ICOND=0; NO AUXILIARY HEATING, 1; OHMIC HEATING, 2; ICRH, 4; ECRH
44 C IIMAX IS NUMBER OF ITERATIONS
45 C PO IS PEAK MICROWAVE POWER IN WATTS
46 C DEA IS INITIAL DENSITY IN 10**9/CC
47 C TWALL IS WALL TEMPERATURE IN EV
48 C AL IS LENGTH IN CM
49 C A IS RADIUS IN CM
50 C AO IS OBSTACLE AREA IN SQ CM
51 C PRES IS NEUTRAL PRESSURE IN 10**5 TORR

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B0 IS FIELD AT OUTER WALL IN KGAUSS
 F IS MICROWAVE FREQUENCY IN GHZ
 G IS THE MICROWAVE CAVITY G
 S IS PUMPING SPEED IN LITERS / SEC
 TMAX IS DURATION OF EXPERIMENT IN SECONDS
 TOFF IS MICROWAVE TURNOFF TIME IN SEC
 NCASE=1
 DQ 800 ICASE=1,NCASE
 ICONDE=1
 IIMAX=1000
 PO=1000.0
 DEAE=0.001
 THALL=0.025
 AL=270.0
 A=18.0
 AO=90.0
 PRES=2.5
 BOE=12.88
 F=2.45
 G=2000.0
 S=100.0
 TMAX=0.005
 TOFF=0.0035

c

SPECIFY INITIAL CONDITIONS

TEASTWALL
 TIASTWALL
 IPRINT=IIMAX/100
 DT=IIMAX/2/FLOAT(IIMAX)
 B=1.0E-4.5
 IIEO
 TIME=0.0
 DDENSE=0.0
 DTIE=0.0
 DTI=0.0
 DNWR=0.0
 DNI=12.88*PRES
 DFC=12.88*PRES
 DNEUT=DNT+DFC
 DEMAX=DEA
 TEMAX=TEA
 TIMAX=TIA
 DNMIN=35.0*PRES
 PRMAX=0.1*PO
 BRMAX=0.1*BO
 SJMAX=0.104*DEA*SQRT(TEA)
 GNMAX=0.0
 SD(2)=IRDI
 SD(3)=IQI
 SD(4)=IFDI
 SD(5)=IGRI
 STI(2)=IC
 STI(3)=EX
 STI(4)=BR
 STI(5)=SR
 STI(6)=TC
 STI(1)=EC
 STI(2)=CX

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109 STI(3)=TCI
110 STI(4)=NCI
111 IOHMIC=MOD(ICOND,2)
112 IHICRH=MOD(ICOND,2,2)
113 IHECRH=MOD(ICOND,4,2)
114 RECORD VALUES
115 WRITE(6,300)
116 300 FORMAT(1H1,'AUXILIARY HEATING:')
117 IF(ICOND.EQ.0) WRITE(6,310)
118 310 FORMAT(1H+,19X,'NONE')
119 IF(IOHMIC.EQ.1) WRITE(6,320)
120 320 FORMAT(1H+,19X,'OHMIC')
121 IF(IHICRH.EQ.1) WRITE(6,330)
122 330 FORMAT(1H+,26X,'ICRH')
123 IF(IHECRH.EQ.1) WRITE(6,340)
124 340 FORMAT(1H+,32X,'ECRH')
125 IF(ICASE.GT.1) WRITE(6,390) ICASE
126 390 FORMAT(1H+,116X,'CASE',I3)
127 WRITE(6,400)
128 400 FORMAT(1H0,2X,'STEP',9X,'TIME',6X,'DENSITY',10X,'TE',11X,'TII',6X,
129 2DNEUTRAL',8X,'FIELD',7X,'POWER',10X,'JSATI',5X,'DETE',11X)
130 WRITE(6,440)
131 440 FORMAT(1H .14X,'(SEC)',3X,'(10**9/CC)',9X,'(EV)',9X,'(EV)',3X,'(10
132 2**9/CC)',5X,'(KGAUSS)',6X,'(WATTS)',4X,'(MA/SQCM)')
133 450 CONTINUE
134 BOLD=B
135 B=B0+SIN(3.14159*TIME/TMAX)*EXP(-0.5*TIME/TMAX)/0.7887+1.0E-5
136 BB=0.34*F/B
137 IF(BR.GT.8.77) G=0.0
138 IF(BR.LE.2.3511) G=0.25744*BR**0.66666667
139 IF(BR.GT.2.3511.AND.BR.LE.8.77) G=2.7*EXP(.45/(BR-8.78))/BR/BR
140 DEAF=0.03*A+DEA/F
141 IF(DEAF.LT.80.0) G=G/EXP(DEAF)
142 IF(DEAF.GE.80.0) G=0.0
143 DEAF=0.06*A+SQRT(AMAX1(DEA-12.35*F*E,0.0))
144 IF(DEAF.LT.80.) G=G+(2.8E-8*DNEUT+9.15E-6*DEA/TEA**1.5)/F/EXP(DEAF)
145 P=PO+AMIN1(1.0,AMAX1(0.0,11.0*TIME*B0**0.6/TMAX-1.7))
146 IF(TIME.GT.TOFF) P=0.0
147 SJI=0.104*DEA*SQRT(AMAX1(TEA,TIA))*(1.0-EXP(-45.0/TEA))
148 DNT0=DNT
149 DFC0=DFC
150 DNT=DNT+DNWR+318.0*8*DT*(322.0*PRES-DNT)/A/A/AL
151 DFC=DFC+0.04*DNWR+318.0*8*DT*(12.80*PRES-DFC)/A/A/AL
152 DNTA=AMIN1(DNT,DNT+(DNT0-DNT)*EXP(-1.2E6*DT*SQRT(TWALL)/A))
153 2*EXP(-4.1E+7*A*D1(DEA,DNEUT,TEA)/DNEUT/SQRT(TWALL))
154 DFCa=AMIN1(DFC,DFC+(DFC0-DFC)*EXP(-3.4E6*DT/A))
155 2*EXP(-1.1E+7*A*D1(DEA,DNEUT,TEA)/DNEUT)
156 DNEUT=DNTA+DFCA
157 GLC=0.0083*SQRT(TEA)+0.28*SQRT(TIA)
158 DNWR=GLC*(D2(DEA,TEA)+D3(DEA,TEA)+D4(DEA)+D5(DEA,TIA))*DT
159 2=0.5*D1(DEA,DNEUT,TEA)+DT
160 DEMAX=AMAX1(DEMAX,DEA)
161 TEMAX=AMAX1(TEMAX,TEA)
162 TIMAX=AMAX1(TIMAX,TIA)
163 DNMIN=AMIN1(DNMIN,DNEUT)
164 BRMAX=AMAX1(BRMAX,B)
165 PPMAX=AMAX1(PPMAX,P)

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166 SUMAX=AMAXI(SUMAX,SI)
167 GNMAY=AMAXI(GNMAY,DEA/DNEUT) GO TO 488
168 IF(MOD(SI+1,IPRINT).NE.1) GO TO 488
169 D(2)=D2(DEA,TEA)
170 D(3)=D3(DEA,TEA)
171 D(4)=D4(DEA)
172 D(5)=D5(DEA,TIA)
173 SAMAXI(0,0,PE2(DEA,TEA,TIA))
174 PE(3)=PE3(DEA,DNEUT,TEA)
175 PE(4)=PE4(DEA,TEA)
176 PE(5)=PE5(DEA,TEA)
177 SAMAXI(0,0,PE6(DEA,TEA))
178 PI(1)=AMAXI(0,0,PI1(2))
179 PI(2)=AMAXI(0,0,PI1(DEA,DNEUT,TIA))
180 PI(3)=AMAXI(0,0,PI16(DEA,TEA,TIA))
181 PI(4)=AMAXI(0,0,PI18(DEA,DNEUT,TIA))
182 DM=0
183 DO 453 I=2,5
184 IF(DM.GT.D(I)) GO TO 453
185 DM=D(I)
186 MD=I
187 CONTINUE
188 EM=0
189 DO 455 I=2,6
190 IF(EM.GT.PE(I)) GO TO 455
191 EM=PE(I)
192 ME=I
193 CONTINUE
194 TIM=0
195 DO 457 I=1,4
196 IF(TIM.GT.PI(I)) GO TO 457
197 TIM=Pi(I)
198 MT=I
199 CONTINUE
200 WRITE(6,460)I,TIME,DEA,TEA,TIA,DNEUT,B,P,SJI,SD(MD),STE(ME),
201 STI(WII)
202 FORMAT(IH,16,8F3.4,5X,3A3)
203 I=I/IPRINT+1
204 T(I)=TIME
205 DES(I)=DEA
206 IE(I)=TEA
207 IIS(I)=TIA
208 DN(I)=DNEUT
209 BR(I)=B
210 PP(SI)=P
211 SJ(SI)=SJI
212 CONTINUE
213 IF(CI.GE.IIMAX) GO TO 640
214 DEAT=DEA
215 TEAT=TEA
216 TIA=TIA
217 DEADEN=DEA+0.5*DENB
218 TEADEN=TEA+0.5*DETE
219 TIAEN=TIA+0.5*DTI
220 DDEN=DI(DEA,DNEUT,TEA)*DT
221 2-D2(DEA,TEA)*DT

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223 3=03(DEA,TEA)*DT
224 4=04(DEA)*DT
225 5=05(DEA,TIA)*DT
226 IF (ABS(DDENS).GT.0.5*DEA) DDENS=SIGN(0.5*DEA,DDENS)
227 DEA=DEAT+0.5*DDENS
228 C INCREMENT ELECTRON TEMPERATURE
229 DUE=PE1(P)*DT
230 2=PE2(DEA,TEA,TIA)*DT
231 3=PE3(DEA,DNEUT,TEA)*DT
232 4=PE4(DEA,TEA)*DT
233 5=PE5(DEA,TEA)*DT
234 6=PE6(DEA,TEA)*DT
235 IF (IOHMIC.EQ.1) DUE=DUE+OHMIC(B,BOLD,DT,TIME,DEA,TEA,DNEUT)*DT
236 IF (IHECRH.EQ.1) DUE=DUE+HECRH(B,TIME,DEA,DNEUT,TEA,A,AL)*DT
237 DTE=(0.66667*DUE+TEA*DDENS)/DEA
238 IF (ABS(DTE).GT.0.5*TEA) DTE=SIGN(0.5*TEA,DTE)
239 TEA=TEAT+0.5*DTE
240 C INCREMENT ION TEMPERATURE
241 DUI=PE2(DEA,TEA,TIA)*DT
242 2=PI3(DEA,DNEUT,TIA)*DT
243 3=PI6(DEA,TEA,TIA)*DT
244 4=PI7(DEA,DNEUT,TEA)*DT
245 5=PI8(DEA,DNEUT,TIA)*DT
246 IF (IHICRH.EQ.1) DUI=DUI+HICRH(B,TIME,DEA,A,AL)*DT
247 DTI=(0.66667*DUI+TIA*DDENS)/DEA
248 IF (ABS(DTI).GT.0.5*TIA) DTI=SIGN(0.5*TIA,DTI)
249 C INCREMENT TIME
250 DEA=DEAT+DDENS
251 TEA=TEAT+DTE
252 TIA=TIAT+DTI
253 II=II+1
254 TIME=DT+FLOAT(II)
255 GO TO 450
256 640 CONTINUE
257 WRITE(6,700) DEMAX,TEMAX,TIMAX,QNMN,BPMAX,PPMAX,SJMAX
258 700 FORMAT(1H0,' MAXIMUM VALUES',7F13.4)
259 FDP=100.0*(DNEUT/334.88-PRES)/PRES
260 WRITE(6,750) FDP
261 750 FORMAT(1H0,14X,' FRACTIONAL CHANGE IN NEUTRAL PRESSURE =',F13.
262 24,' PER CENT')
263 WRITE(6,760) QNMN
264 760 FORMAT(1H0,28X,' MAXIMUM RATIO OF NE / NO =',F13.4)
265 800 CONTINUE
266 C GRAPH OUTPUT
267 DO 820, I=1,101
268 T(I)=6.0*T(I)/TMAX
269 DF(I)=6.0*DE(I)/DEMAX
270 TE(I)=6.0*TE(I)/TEMAX
271 TI(I)=6.0*TI(I)/TIMAX
272 DN(I)=6.0*DN(I)/DNEUT
273 BP(I)=6.0*BP(I)/BPMAX
274 PP(I)=6.0*PP(I)/PPMAX
275 820 SJ(I)=6.0*SJ(I)/SJMAX
276 CALL GRAPH2(T,'R',DE,'R',101,16X61,'NONE', 'MULTIPOLE SIMULATION',
277 2,' TIME', ' DENSITY AND TEMP (NORMALIZED)', 'D')
278 CALL GRPH2V(T,'R',TE,'R',101,'NONE', 'E')
279 CALL GRPH2V(T,'R',TI,'R',101,'NONE', 'I')

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280 CALL GRPH2V(T,IR, DN,IR,101,INONE,INI)
281 CALL GRAPH2(T,IR, BP,IR,101,16X6,INONE),MULTIPOLE SIMULATION.,I
282 2,ITIME,1,IFIELD AND POWER (NORMALIZED),I,(BI)
283 CALL GRPH2V(T,IR,PP,IR,101,INONE,PI)
284 CALL GRPH2V(T,IR,SJ,IR,101,INONE,IJI)
285 CALL GRPHND
286 END

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END OF COMPILATION: NO DIAGNOSTICS.

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S ,OHMIC
N=MACC 1.178-06/08/76-14:44:19 OHMIC
1 FUNCTION OHMIC(B,BOLD,DT,T,DENS,TE,DNEUT)
2 BT(T)=0.177*BTMAX*SIN(1571.0*(T-TON))/EXP(654.0*(T-TON))
3 OHMIC=0.0
4 TON=0.0035
5 IF(T.LE.TON) RETURN
6 BTMAX=1.25
7 ZEFF=1.0
8 DBPO=0.1*(B-BOLD)/DT
9 DBT=0.208*(BT(T+0.5*DT)-BT(T-0.5*DT))/DT
10 ALF=0.141*B/BT(T)
11 OHMIC=(6.51*ALF**2*DBT**2/(0.591+ALF**2)-12.6*ALF*DBT*DBPO/(0.663+
12 2*ALF**2)+8.79*DBPO**2/(0.962+ALF**2))
13 3/(ZEFF/1600.0/TE**1.5+6.29E-6*DNEUT/DENS+0.025*SQRT(TE)/DENS)
14 OHMIC=MAXI(OHMIC,0.0)
15 RETURN
16 END

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END OF COMPILATION: NO DIAGNOSTICS.

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S ,HICRH
N=MACC 1.178-06/08/76-14:44:30 HICRH
1 FUNCTION HICRH(B,TIME,DEA,A,AL)
2 C SEE BK 6 PG. 118 FOR EXPLANATIONS AND REFERENCES
3 HICRH=0.0
4 TON=0.0045
5 TOFF=0.006
6 IF(TIME.LT.TON.OR.TIME.GT.TOFF) RETURN
7 VQPO=6.5
8 RPAR=800.0
9 RO=200
10 FREQ=2.0
11 DRDT=2.0E4
12 TRISE=5.E-5
13 BO=.62A*FREQ/B
14 IF(BO.LT.0.51198402) GI=.027*BO**1.03
15 IF(BO.GE.0.51198402.AND.BO.LT.0.82973281) GI=.051*BO**1.98
16 IF(BO.GE.0.82973281.AND.BO.LT.1.01377) GI=.019/BO**3.31
17 IF(BO.GE.1.01377) GI=.0185/BO**1.36
18 RLOAD=1.32E4*FREQ*AL/DEA/A/A/GI
19 REQ=RPAR+RLOAD/(RPAR+RLOAD)
20 VOP=VQPO*(1.87*REQ+RO)/(1.87*REQ+RO+DRDT*(TIME-TON))
21 XP=(TIME-TON)/TRISE
22 IF(XP.LT.88.) VOP=VOP*(1.0-EXP(-XP))
23 HICRH=7.62E10*DEA*VOP*VOP*GI/AL/AL/FREQ
24 RETURN
25 END

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END OF COMPILATION; NO DIAGNOSTICS.

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9. HECRH
N=MACC 1.178-06/08/76-14:44:41 HECRH
1 FUNCTION HECRH(B,TIME,DEA,DNEUT,TEA,A,AL)
2 TON=0.003
3 TOFF=0.004
4 HECRH=0.0
5 IF(TIME.LT.TON.OR.TIME.GT.TOFF) RETURN
6 F=9.0
7 DEAF=0.03*A+DEA/P
8 P=2.0E4
9 Q=3000.0
10 BO=0.34+F/B
11 IF(BO.GT.8.77) G=0.0
12 IF(BO.LE.2.3511) G=0.25744*BO**0.66666667
13 IF(BO.GT.2.3511.AND.BO.LE.8.77) G=2.7*EXP(.45/(BO-8.78))/BO/BO
14 IF(DEAF.LT.80.0) G=G/EXP(DEAF)
15 IF(DEAF.GE.80.0) G=0.0
16 DEAF=0.06*A+SQRT(AMAX1(DEA-12.35*F/F,0.0))
17 IF(DEAF.LT.80.)G=G+(2.8E-8*DNEUT+9.15E-6*DEA/TEA**1.5)/F/EXP(DEAF)
18 HECRH=2.0E9*P*G*DEA/(G*DEA+12.35*F*F/Q)/A/A/AL
19 RETURN
20 END

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END OF COMPILATION; NO DIAGNOSTICS.

XN
4 RL1B64 06/08-14:45:09
P

AUXILIARY HEATING: OHMIC

STEP	TIME (SEC)	DENSITY (10 ²⁰ /CC)	TE (eV)	TI (eV)	DNEUTRAL (10 ²⁰ /CC)
0	.0000	.0010	.0250	.0250	837.2000
10	.0000	.0005	.0375	.0370	837.2040
20	.0001	.0005	.0376	.0366	837.2040
30	.0001	.0005	.0377	.0361	837.2040
40	.0002	.0005	.0377	.0357	837.2040
50	.0002	.0005	.0378	.0353	837.2040
60	.0003	.0005	.0378	.0349	837.2040
70	.0003	.0005	.0379	.0349	837.2040
80	.0004	.0005	.0380	.0349	837.2040
90	.0004	.0005	.0380	.0338	837.2040
100	.0005	.0005	.0381	.0339	837.2040
110	.0005	.0005	.0381	.0332	837.2040
120	.0006	.0005	.0382	.0329	837.2040
130	.0006	.0005	.0382	.0326	837.2040
140	.0007	.0005	.0383	.0325	837.2040
150	.0007	.0005	.0383	.0321	837.2040
160	.0008	.0005	.0384	.0318	837.2040
170	.0008	.0005	.1453	.0318	837.2040
180	.0009	.0005	.3014	.0316	837.2036
190	.0009	.0050	.1626	.0277	837.1897
200	.0010	.0296	.7402	.0279	837.0671
210	.0010	.5171	.35417	.0279	835.6402
220	.0011	.9511	.2284	.0272	819.7187
230	.0011	.6655	.7768	.0329	815.8335
240	.0012	.6215	.7719	.0466	818.2517
250	.0012	.0195	.5981	.0649	818.1283
260	.0013	.3146	.1550	.0844	817.5303
270	.0013	.3253	.8914	.1039	816.9216
280	.0014	.3303	.6215	.1234	816.4829
290	.0014	.3322	.4306	.1428	816.0901
300	.0015	.1858	.3017	.1617	815.7215
310	.0015	.4475	.2118	.1796	815.3743
320	.0016	.6394	.1467	.1972	815.0459
330	.0016	.7818	.0983	.2137	814.7334
340	.0017	.8873	.0617	.2292	814.4342
350	.0017	.9680	.0338	.2438	814.1459
360	.0018	.0215	.0125	.2575	813.8666
370	.0018	.0618	.9962	.2703	813.5947
380	.0019	.0895	.9838	.2822	813.3288
390	.0019	.1076	.9746	.2933	813.0678
400	.0020	.1184	.9679	.3036	812.8108
410	.0020	.1238	.9633	.3132	812.5569
420	.0021	.1252	.9603	.3221	812.3055
430	.0021	.1239	.9587	.3303	812.0560
440	.0022	.1208	.9582	.3379	811.8079
450	.0022	.1169	.9586	.3449	811.5608
460	.0023	.1054	.9603	.3513	811.3154
470	.0023	.0779	.9652	.3572	811.0715
480	.0024	.0375	.9736	.3626	810.8271
490	.0024	.9868	.9852	.3675	810.5812
500	.0025	.9283	.9994	.3718	810.3326
510	.0025	.8642	.0158	.3755	810.0805
520	.0026	.7963	.0341	.3788	809.8240

B_p	μ wave	JSAT	dominating term
FIELD (KGAUSS)	POWER (WATTS)	(MA/SQCM)	DE TE TI
0000	0000	0000	RD BR NC
0349	0000	0000	OL TC NC
0694	0000	0000	OL TC NC
1034	0000	0000	OL TC NC
1371	0000	0000	OL TC NC
1702	0000	0000	OL TC NC
2029	0000	0000	OL TC NC
2350	0000	0000	OL TC NC
2666	0000	0000	OL TC NC
2976	0000	0000	OL TC NC
3280	0000	0000	OL TC NC
3577	0000	0000	OL TC NC
3868	0000	0000	OL TC NC
4152	0000	0000	OL TC NC
4430	0000	0000	OL TC NC
4700	0000	0000	OL TC NC
4962	0000	0000	OL TC NC
5217	31, 9335	0000	OL TC NC
5464	133, 8119	0002	OL EX TC
5703	235, 6903	0011	OL TC TC
5934	337, 5688	0092	OL TC TC
6157	439, 4472	1280	OL TC TC
6371	541, 3257	3, 0583	OL EX TC
6577	643, 2041	7, 4548	OL EX TC
6774	745, 0825	7, 4657	OL EX NC
6963	846, 9610	7, 3844	OL EX NC
7142	948, 8394	7, 4293	OL EX NC
7313	1000, 0000	7, 4835	OL EX NC
7474	1000, 0000	7, 4454	OL EX NC
7626	1000, 0000	7, 4038	OL EX NC
7769	1000, 0000	7, 3737	OL EX NC
7903	1000, 0000	7, 3521	OL EX NC
8028	1000, 0000	7, 3358	OL EX NC
8143	1000, 0000	7, 3228	OL EX NC
8249	1000, 0000	7, 3122	OL EX NC
8346	1000, 0000	7, 3033	OL EX NC
8433	1000, 0000	7, 2957	OL EX NC
8511	1000, 0000	7, 2893	OL EX NC
8579	1000, 0000	7, 2838	OL EX NC
8638	1000, 0000	7, 2791	OL EX NC
8688	1000, 0000	7, 2752	OL EX NC
8729	1000, 0000	7, 2721	OL EX NC
8760	1000, 0000	7, 2696	OL EX NC
8782	1000, 0000	7, 2679	OL EX NC
8796	1000, 0000	7, 2668	OL EX NC
8800	1000, 0000	7, 2663	OL EX NC
8795	1000, 0000	7, 2655	OL EX NC
8782	1000, 0000	7, 2643	OL EX NC
8760	1000, 0000	7, 2636	OL EX NC
8729	1000, 0000	7, 2636	OL EX NC
8690	1000, 0000	7, 2644	OL EX NC
8642	1000, 0000	7, 2660	OL EX NC
8586	1000, 0000	7, 2683	OL EX NC

STEP	TIME (SEC)	Ne ³⁵⁻ (10 ⁹ /cm ³)	Te (eV)	Ti (eV)	neutral density
30	,0026	34,7261	0540	,3815	809,5625
40	,0027	34,6550	0752	,3835	809,2953
50	,0027	34,5841	0975	,3856	809,0218
60	,0028	34,5144	1207	,3870	808,7415
70	,0028	34,4466	1445	,3880	808,4540
80	,0029	34,3813	1689	,3886	808,1589
90	,0029	34,3192	1937	,3888	807,8559
100	,0030	34,2604	2189	,3887	807,5446
110	,0030	34,2053	2444	,3883	807,2247
120	,0031	34,1541	2701	,3876	806,8960
130	,0031	34,1068	2960	,3866	806,5581
140	,0032	34,0635	3222	,3854	806,2108
150	,0032	34,0243	3486	,3840	805,8537
160	,0033	33,9890	3753	,3823	805,4867
170	,0033	33,9576	4022	,3805	805,1093
180	,0034	33,9300	4296	,3785	804,7211
190	,0034	33,9061	4574	,3763	804,3218
200	,0035	33,8858	4857	,3740	803,9109
210	,0035	33,8739	5146	,3714	803,4884
220	,0036	251,4967	5421	,1741	677,1127
230	,0036	484,4878	5702	,1050	411,6204
240	,0037	644,5074	6027	,1555	179,3079
250	,0037	732,2671	6353	,2547	89,9616
260	,0038	778,2188	6679	,3701	58,7936
270	,0038	800,6711	7003	,4872	45,6386
280	,0039	809,8063	7329	,6003	39,2032
290	,0039	811,5023	7651	,7074	35,8304
300	,0040	809,2790	7969	,8082	34,0830
310	,0040	805,2587	8289	,9030	33,3086
320	,0041	800,7044	8606	,9929	33,1867
330	,0041	796,3375	8922	,1,0788	33,5598
340	,0042	792,5302	9238	,1,1620	34,3638
350	,0042	789,4208	9550	,1,2433	35,5962
360	,0043	786,9750	9864	,1,3244	37,3026
370	,0043	785,0103	10176	,1,4058	39,5739
380	,0044	783,1856	10483	,1,4884	42,5507
390	,0044	780,9617	10784	,1,5729	46,4337
400	,0045	777,5336	11079	,1,6598	51,5009
410	,0045	771,7416	11368	,1,7496	58,1274
420	,0046	761,9807	11653	,1,8422	66,8034
430	,0046	745,7857	11934	,1,9370	78,1372
440	,0047	715,5908	12211	,2,0269	92,7370
450	,0047	645,0611	12484	,2,0547	112,8447
460	,0048	507,0191	12753	,1,8504	148,3165
470	,0048	327,0318	13018	,1,3639	212,4293
480	,0049	162,7962	13279	,7998	289,4035
490	,0049	49,7327	13536	,3499	346,5477
1000	,0050	1,3780	8,0308	,0865	376,8779
				,0802	390,5984

MAXIMUM VALUES → 811,5647 586,7402 2,0611 33,1663

FRACTIONAL CHANGE IN NEUTRAL PRESSURE → -53,3447%

MAXIMUM RATIO OF NE / NO → 24,3365

B_{po} (kG)	μ wave Power (watts)	probe ion sat. current ($\mu A/cm^2$)	dominating term $N_e \ T_e \ T_i$
8522	1000,0000	7,2735	OL EX NC
8450	1000,0000	7,2756	OL EX NC
8371	1000,0000	7,2805	OL EX NC
8283	1000,0000	7,2863	OL EX NC
8189	1000,0000	7,2930	OL EX NC
8087	1000,0000	7,3006	OL EX NC
7977	1000,0000	7,3090	OL EX NC
7861	1000,0000	7,3184	OL EX NC
7738	1000,0000	7,3286	OL EX NC
7609	1000,0000	7,3398	OL EX NC
7473	1000,0000	7,3518	OL EX NC
7331	1000,0000	7,3648	OL EX NC
7183	1000,0000	7,3788	OL EX NC
7029	1000,0000	7,3937	OL EX NC
6870	1000,0000	7,4095	OL EX NC
6705	1000,0000	7,4265	OL EX NC
6536	1000,0000	7,4445	OL EX TC
6361	1000,0000	7,4635	OL EX TC
6182	,0000	34,6904	OL EX TC
5998	,0000	111,9706	OL EX TC
5810	,0000	213,0038	OL EX TC
5618	,0000	274,5396	OL TC TC
5423	,0000	303,4991	OL TC TC
5223	,0000	317,4632	OL TC TC
5021	,0000	324,8130	OL TC TC
4815	,0000	329,2412	OL TC TC
4607	,0000	332,4708	OL TC TC
4396	,0000	335,3200	OL TC TC
4183	,0000	338,1810	OL TC TC
3968	,0000	341,0186	OL TC TC
3751	,0000	343,8681	OL TC TC
3532	,0000	346,4847	OL TC TC
3312	,0000	348,5667	OL TC TC
3090	,0000	349,7242	OL TC TC
2868	,0000	349,4802	OL TC TC
2645	,0000	347,2707	OL TC TC
2422	,0000	342,4558	OL TC TC
2199	,0000	334,3605	OL TC TC
1975	,0000	322,3689	PD EX TC
1752	,0000	306,0895	PD EX TC
1529	,0000	285,4403	PD EX TC
1307	,0000	259,1797	PD EX TC
1086	,0000	219,8228	GR EX TC
0865	,0000	161,0594	GR TC TC
0647	,0000	97,0481	GR TC TC
0429	,0000	46,7271	GR TC TC
0214	,0000	14,3678	GR TC TC
.0000	.0000	.2042	RD TC TC
8800	1000,0000	349,8245	← MAX VALUES

PER CENT

OFFIN

RUNID: C00295 PROJECT: 02980

USER: 4126810219

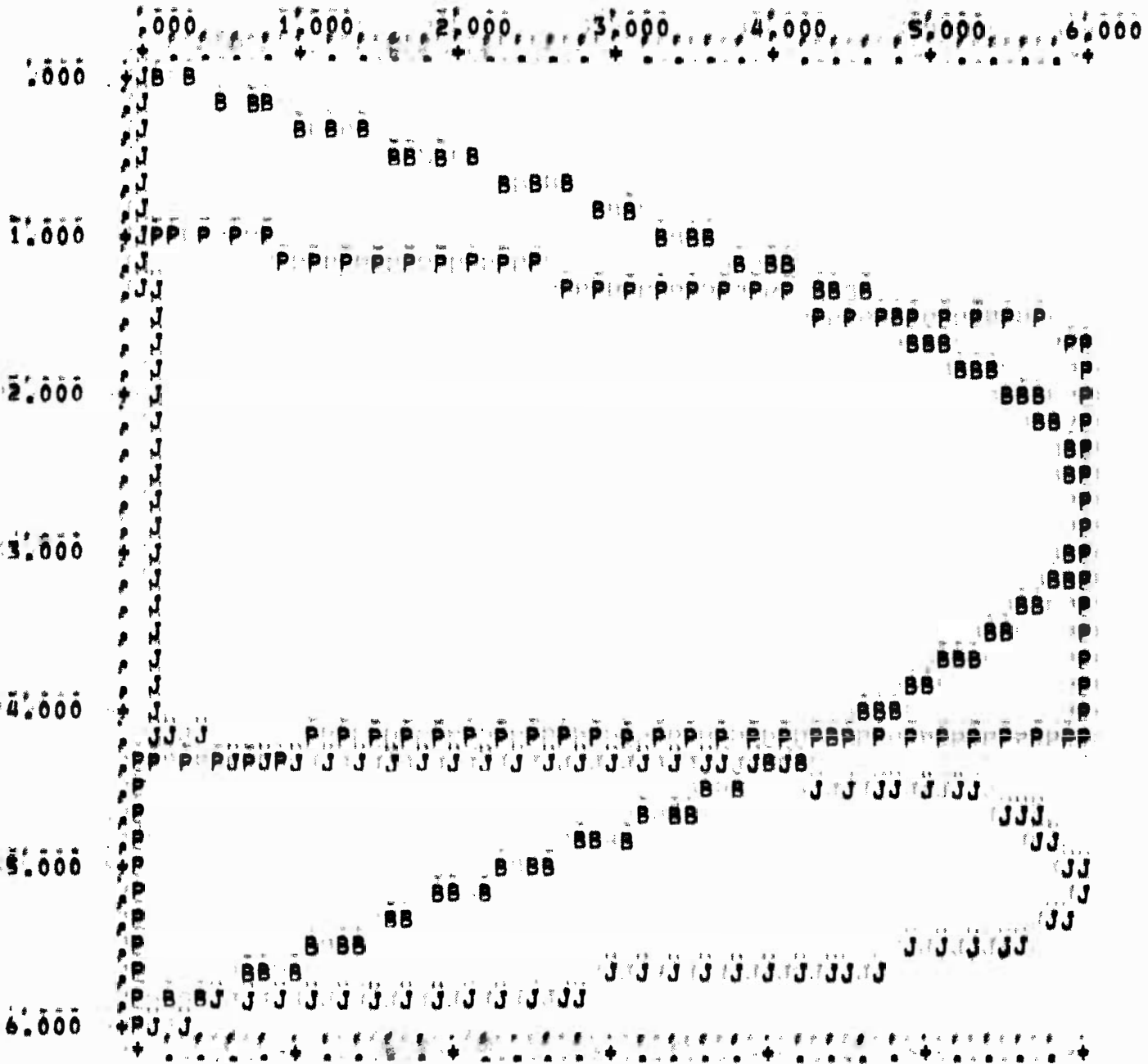
ITEM	AMOUNT	COST (DOLLARS)
CPU TIME	00:00:14.047	\$0.53
FILE I/O REQUESTS	558	\$0.25
FILE I/O WORDS	471145	\$0.21
SWAP REQUESTS	1	\$0.00
SWAP WORDS	3072	\$0.00
MEMORY USAGE	0.383	\$0.24
CARDS IN	355	\$0.09
PAGES PRINTED	13	\$0.18
ER + CC	12	\$0.12
JOB CHARGE	1	\$0.05
TOTAL COST		\$1.67

THE ABOVE DOLLAR AMOUNTS ARE APPROXIMATE AND ARE BASED ON RATES FOR STANDARD RUNS
USER BALANCE \$113.46

INITIATION TIME: 14:43:15 JUN 8, 1976
TERMINATION TIME: 14:49:21 JUN 8, 1976
PREVIOUS RUN TIME: 14:13:59 MAY 20, 1976

MULTIPOLE SIMULATION

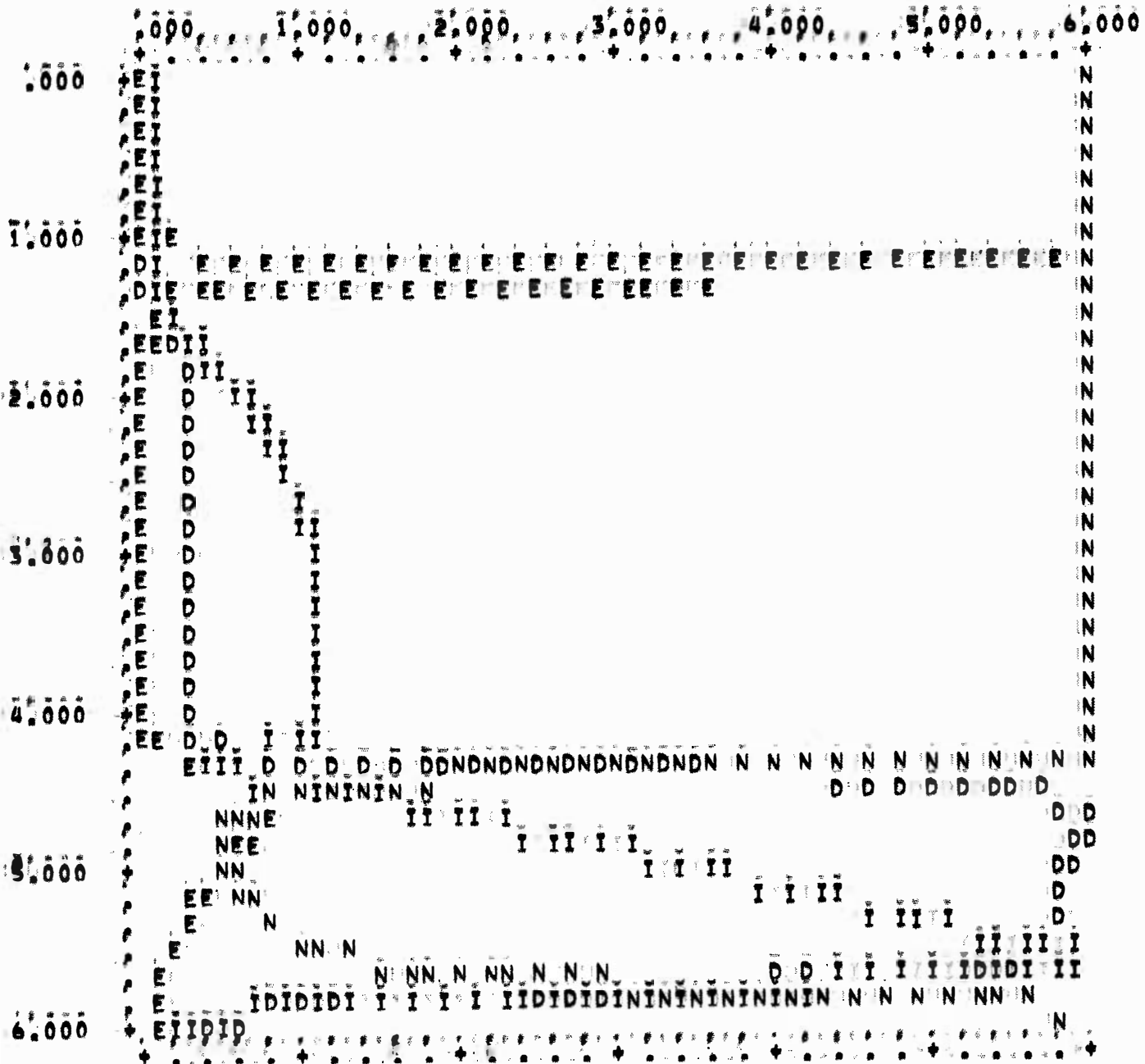
FIELD AND POWER (NORMALIZED)



TIME (normalized)

MULTIPOLE SIMULATION

DENSITY AND TEMP (NORMALIZED)



TIME (normalized)