

POSSIBILITIES FOR PROLONGING THE TOKAPOLE II DISCHARGE

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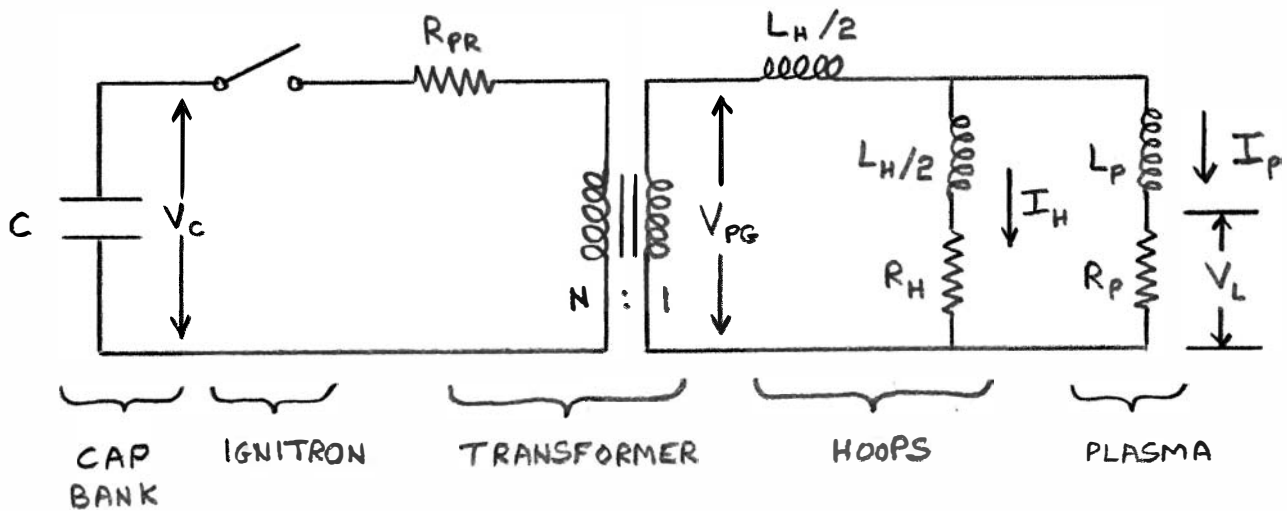
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POSSIBILITIES FOR PROLONGING THE TOKAPOLE II DISCHARGE

J.C. Sprott

It is desirable to extend the duration of the Tokapole II discharge current pulse for a number of reasons: (1) to facilitate the study of slowly growing instabilities, (2) to insure that equilibrium is achieved with wall reflxed neutrals, and (3) to achieve higher temperatures and poloidal beta. In this note, an extension of the circuit model of PLP 756 will be used to determine what changes to the external poloidal field circuitry would be required to prolong the plasma current pulse in Tokapole II.

The circuit model used is as follows



The values assumed for the circuit components are:

Capacitance of B_p bank: $C = 0.0072 \text{ F}$

Primary resistance: $R_{PR} = 0.0485 \Omega$

Turns ratio: $N = 40$

Hoop inductance: $L_H = 0.219 + 0.06834 (1 - e^{-396t}) \mu\text{H}$

$$\text{Hoop resistance: } R_H = 7.8 + 16.9 e^{-396t} \mu\Omega$$

$$\text{Plasma inductance: } L_P = 0.8 - 0.5 L_H \mu\text{H}$$

$$\text{Plasma resistance: } R_P = 12.26 Z_{\text{eff}}/a^2 T_e^{3/2} \text{ ohms}$$

The primary resistance and hoop inductance at $t = 0$ were taken from PLP 744. The hoop resistance and inductance correction are approximations to the calculation of Spencer (PLP 771 and subsequent private communications), where t is the time in seconds after the ignition fires. The plasma inductance was measured by Miller (PLP 756). The plasma resistance comes from Spitzer, where Z_{eff} is taken equal to one, a is the distance in cm from the geometric axis to the field nulls, and T_e is the plasma electron temperature in eV. The quantity a is determined from the ratio of plasma current I_P to hoop current I_H by

$$a = 21 (I_P/I_H)^{1/4} \text{ cm}$$

The electrical circuit equations which relate the variables, V_C , I_H , and I_P are:

$$C \frac{dV_C}{dt} = - \frac{I_H + I_P}{N}$$

$$\frac{L_H}{2} \frac{d(I_H + I_P)}{dt} + \frac{L_H}{2} \frac{dI_H}{dt} + R_H I_H = V_{PG}$$

$$\frac{L_H}{2} \frac{d(I_H + I_P)}{dt} + L_P \frac{dI_P}{dt} + R_P I_P = V_{PG}$$

where

$$V_{PG} = \frac{V_C}{N} - \frac{R_{PR} (I_H + I_P)}{N^2}$$

With a bit of algebra, these equations can be put in a more convenient form:

$$\frac{dV_C}{dt} = - \frac{I_H + I_P}{NC} \quad (1)$$

$$\frac{dI_H}{dt} = \frac{2V_{PG} L_P/L_H - R_H (2L_P/L_H + 1) I_H + R_P I_P}{2L_P + \frac{1}{2} L_H} \quad (2)$$

$$\frac{dI_P}{dt} = \frac{V_{PG} + R_H I_H - 2R_P I_P}{2L_P + \frac{1}{2} L_H} \quad (3)$$

The set of non-linear differential equations (1)-(3) can in principle be solved provided $T_e(t)$ is known. Since $T_e(t)$ is not known a priori, it is better to calculate it self-consistently from the ohmic heating input power and some reasonable model of the energy loss rate:

$$\frac{d}{dt} (1.5 n_e T_e v) = I_P^2 R_P - 1.5 n_e T_e v / \tau_E \quad (4)$$

In equation (4) the density n is taken as a constant equal to $1 \times 10^{13} \text{ cm}^{-3}$, and the volume v is taken as $6 \times 10^5 \text{ cm}^3$. The energy confinement time τ_E is taken as an adjustable parameter whose value is time-independent. These assumptions are reasonable because tokamak densities tend to be determined only by the toroidal field strength, and energy confinement times are determined almost entirely by the density (see PLP 753). Late gas puffing may be

required to keep the density constant throughout the discharge. Note also that the loss term in equation (4) contains the quantity n/τ_E which presumably remains approximately constant even if n varies somewhat.

The set of equations (1)-(4) has been solved numerically using program TOKII listed in Appendix A. The initial conditions are normally taken as

$$V_C = 2000 \text{ volts}$$

$$I_H = 0$$

$$I_P = 0$$

$$T_e = 5 \text{ eV}$$

The resulting plasma current as a function of time for various values of τ_E is shown in figure 1. Also shown is the result of a typical experimental pulse with the divertor baffle inserted to reduce the current which flows in the common flux region. The experiment is in reasonable agreement with the model provided the energy confinement time is the order of a few hundred μsec . The discrepancy may result from some combination of the following: (1) an increase in the plasma inductance as a function of time, (2) a decrease in energy confinement time as a function of time, (3) a rise in plasma density as a function of time, (4) a rise in Z_{eff} as a function of time, or (5) a shift in the current distribution toward the rings as a function of time. Note that the current waveform is a sensitive function of confinement time for τ_E in the range below about 400 μsec , but for high τ_E , the plasma inductance dominates, and the plasma and hoops form an inductive current divider so that the plasma current follows the hoop current according to

$$I_P = L_H I_H / 2L_P \approx 0.20 I_H$$

This result is further exhibited in figure 2 where the amp-seconds ($\int I_p dt$) are plotted as a function of τ_E .

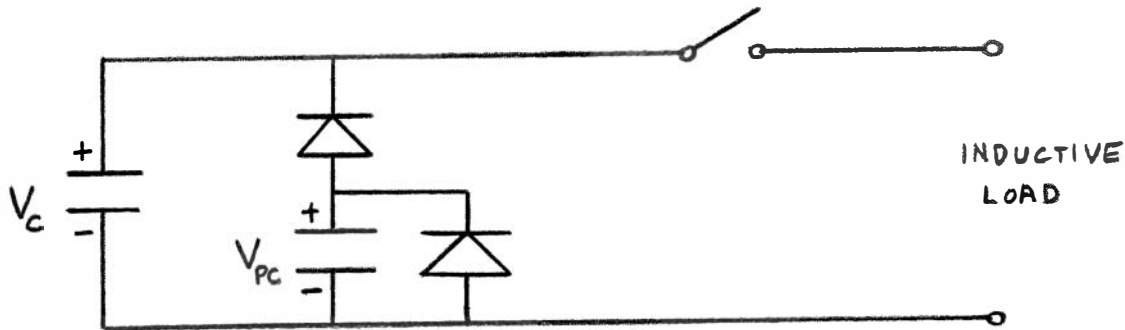
The amp-seconds is apparently a good figure of merit for low τ_E , but not for high τ_E . As a further point of interest, the electron temperature as a function of time is plotted in figure 3 for various values of assumed energy confinement time. Note that the electron temperature is volume-averaged over the entire machine. The central current channel temperature is undoubtedly higher, but then so probably is Z_{eff} , and so the current waveforms are probably reasonable.

In order to extend the discharge time, several solutions come to mind. Perhaps the simplest is to increase the primary resistance and raise the capacitor voltage to compensate for the additional IR drop. Figure 4 shows such a case with $V_C = 5000$ volts and $R_{PR} = 0.5 \Omega$. Although the high initial gap voltage ($V_{PG} = 125$ volts) does help the discharge get started, especially at low τ_E , the decay of the plasma current is disappointingly rapid. The $\tau_E = 1600 \mu\text{sec}$ case still only has 130 amp-seconds and a peak electron temperature of 98 eV.

A second possible solution is to passively crowbar the poloidal field so that the capacitor voltage is clamped at zero rather than allowed to swing negative. The result is shown in figure 5. Little improvement over the uncrowbarred case (Fig. 1) is noted. The $\tau_E = 1600 \mu\text{sec}$ case has 139 amp-seconds and a peak electron temperature of 89 eV.

The next level of complexity would be a power crowbar which will hold the capacitor bank voltage constant at a prescribed voltage (V_{PC}) whenever it attempts to drop below that voltage. In practice this is done by switching a large bank of capacitors charged to V_{PC} in parallel with the initial capacitor

bank at the prescribed time. Actually a diode would in principle suffice as shown below:



A second diode is also shown in parallel with the power crowbar bank to provide a passive crowbar when the voltage eventually tries to reverse. This extra diode is also desired to prevent the power crowbar bank voltage from reversing in the event that electrolytic capacitors are used. We begin by examining the case of an ideal power crowbar (that is, one with an infinite capacitance) and return later to the question of how large a bank is actually required. Figure 6 shows the result of such a case with $V_{PC} = 450$ volts. This value of voltage was chosen because it is the highest voltage in which electrolytic capacitors are commonly available. It is clear that a spectacular improvement is achieved for $\tau_E \geq 1000 \mu\text{sec}$. The $1600 \mu\text{sec}$ case has 288 amp-seconds (ignoring times later than 10 msec) and a peak electron temperature of 94 eV. At 10 msec the loop voltage V_L and the poloidal gap voltage V_{PG} are about 2 volts and the core flux is about 0.1 volt-sec which is about the limit without cocking.

Another case was examined in which the power crowbar voltage was 900 volts (requiring a series parallel combination of electrolytics), and the result is shown in figure 7. This is a case of overkill. The $1600 \mu\text{sec}$ case has 491 amp-seconds (up to 10 msec) and a peak electron temperature of 174 eV. At 10 msec the loop voltage V_L is 2 volts and the poloidal gap voltage V_{PG} is 7 volts. The core flux is 0.15 volt-sec which might barely be achievable with an elaborate core cocking circuit.

Now we return to the question of how much capacitance is required in a reasonable power crowbar bank. We take as a standard the $\tau_E = 800 \mu\text{sec}$ case in figure 6 and vary the capacitance C_{PC} . The result is shown in figure 8. To take full advantage of the 10 msec pulse apparently requires $C_{PC} \gtrsim 0.5$ farads.

The variation of $\int I_p dt$ vs τ_E for the crowbarred case with $V_{PC} = 450$ volts and $C_{PC} = 0.4$ farads is shown in figure 9. Comparison with figure 2 shows that the amp-seconds is a reasonable figure of merit in the presence of a suitable power crowbar, although it becomes less sensitive as τ_E increases.

In summary, it has been shown that in the absence of a power crowbar, the plasma current waveform can never exceed in duration that of the hoop current, or in amplitude a value of about 20% of the hoop current. With a power crowbar, the amplitude is still limited to about 20% of the hoop current, but the duration is limited only by saturation of the iron core. A power crowbar of ~ 0.5 F and 450 volts should be about right for extending the current pulse to ~ 10 msec without saturating the core, assuming an energy confinement time of ~ 1 msec.

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KII
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APPENDIX A

TOKII

```
C PROGRAM TOKII = J. C. SPROTT = JULY 2, 1978
  DIMENSION YI(4),YF(5),SAVE(52,6)
  COMMON C,RPR,N,ALP,DENS,ZEFF,TAUE
  EXTERNAL DERIVS
  DO 900 ICASE=1,9
  C=0.0072
  RPR=0.0485
  N=40
  VC=2000.0
  ZEFF=1.0
  TAUE=1.0E-4*1.4142136**(ICASE-1)
  AIH=0.1
  AIP=0.1
  DENS=1.0E13
  TE=3.0
  AMPSEC=0.0
  VSEC=0.0
  YI(1)=VC
  YI(2)=AIH
  YI(3)=AIP
  YI(4)=1.44E-13*DENS*TE
  WRITE(6,300)
300 FORMAT(1H1,' TIME IP VPG IHoop VLOOP VCAP P
20H TE')
  CALL DEPC(4,0.0,YI,0.010,YF,DERIVS,1.0E-3,0.0,1.0E-7,1.0E-3,1.0E-4
2,2,2.0E-4,2H.,SAVE,52,6,NPOINT,NOTIFY,S400)
400 CONTINUE
  DO 600 I=1,NPOINT
  IT=200*(I-1)
  IP=1.0E-3*SAVE(I,3)
  IVPG=SAVE(I,5)
  IH=1.0E-3*SAVE(I,2)
  IVL=SAVE(I,4)
  IVC=SAVE(I,1)
  IPOH=1.0E-3*SAVE(I,3)*SAVE(I,4)
  ITE=SAVE(I,6)
  WRITE(6,500) IT,IP,IVPG,IH,IVL,IVC,IPOH,ITE
500 FORMAT(1H ,8I8)
  AMPSEC=AMPSEC+2.0E-4*SAVE(I,3)
  VSEC=VSEC+2.0E-4*AMAX1(SAVE(I,5),0.0)
600 CONTINUE
  WRITE(6,700) AMPSEC
700 FORMAT(1H0,'AMP SECONDS',F11.6)
  WRITE(6,720) VSEC
720 FORMAT(1H0,'VOLT SECONDS',F10.6)
900 CONTINUE
  END
```

COMPILATION: NO DIAGNOSTICS.
IVS

1.17S-01/05/79-14:37:02

DERIVS

```

SUBROUTINE DERIVS(T,Y,DY,STORE,ITEST)
DIMENSION Y(1),DY(1),STORE(1)
COMMON C,RP,N,ALP,DENS,ZEFF,TAUE
VC=Y(1)
AIH=Y(2)
AIP=Y(3)
CO=C
VCROW=450.0
CCROW=0.5
IF(VC.LE.VCROW) CO=C+1.0E4*C*(VCROW-VC)
IF(CO.GT.CCROW) CO=CCROW
IF(VC.LE.50.0) CO=CCROW+1.0E4*CCROW*(50.0-VC)
RH=7.8E-6+16.9E-6*EXP(-396.0*T)
ALH=0.319E-6+0.06834E-6*(1.0-EXP(-396.0*T))
ALP=0.8E-6-0.5*ALH
TE=Y(4)/DENS/1.44E-13
A=AMAX1(3.0,AMIN1(31.11,21.0*(ABS(AIP/AIH))*0.25))
RP=12.26*ZEFF/A/A/(AMAX1(TE,0.025))*1.5
VL=AIP*RP
VPG=VC/FLOAT(N)-RPR*(AIH+AIP)/FLOAT(N*N)
DY(1)=- (AIH+AIP)/CO/FLOAT(N)
DY(2)=(2.0*VPG*ALP/ALH-RH*(2.0*ALP/ALH+1.0)*AIH+RP*AIP)
2/(2.0*ALP+0.5*ALH)
DY(3)=(VPG+RH*AIH-2.0*RP*AIP)/(2.0*ALP+0.5*ALH)
DY(4)=AIP*VL-1.44E-13*DENS*TE/TAUE
DO 300 I=1,3
300 STORE(I)=Y(I)
STORE(4)=VL
STORE(5)=VPG
STORE(6)=TE
RETURN
END

```

COMPILATION: NO DIAGNOSTICS.

\$ 01/05-14:37:10

TIME	IP	VPG	IHOOP	VLOOP	VCAP	POH	TE
0	0	49	0	0	2000	0	5
200	3	48	42	15	1983	52	7
400	6	45	81	12	1937	73	13
600	9	42	115	9	1863	90	20
800	13	39	146	8	1764	104	28
1000	16	35	173	6	1642	114	35
1200	19	30	196	5	1501	122	42
1400	22	26	215	5	1343	128	48
1600	24	21	230	5	1172	131	53
1800	26	16	241	5	990	132	58
2000	27	11	249	4	800	131	61
2200	27	6	252	4	607	128	63
2400	27	2	253	4	446	123	65
2600	27	2	252	4	443	119	65
2800	27	2	252	4	440	117	65
3000	26	2	252	4	437	115	65
3200	26	2	251	4	435	113	65
3400	25	2	251	4	432	112	64
3600	25	2	251	4	429	110	63
3800	25	2	250	4	426	109	63
4000	24	2	250	4	424	107	62
4200	23	2	250	4	421	106	61
4400	23	2	250	4	418	104	61
4600	22	2	250	4	415	102	60
4800	22	2	250	4	413	100	59
5000	21	2	250	4	410	98	58
5200	21	1	250	4	407	96	57
5400	20	1	250	4	404	94	56
5600	19	1	250	4	402	92	55
5800	18	1	250	4	399	90	54
6000	18	1	250	4	396	87	53
6200	17	1	250	4	394	85	52
6400	16	1	250	4	391	82	50
6600	15	1	250	5	388	79	49
6800	15	1	250	5	386	76	48
7000	14	1	250	5	383	73	46
7200	13	1	250	5	380	69	45
7400	12	1	250	5	378	65	43
7600	11	1	250	5	375	61	41
7800	10	1	250	5	373	57	39
8000	9	1	250	5	370	52	37
8200	8	1	250	5	367	47	35
8400	6	1	250	5	365	41	33
8600	5	1	250	6	362	35	30
8800	4	1	250	6	360	28	27
9000	3	1	250	6	357	20	24
9200	1	1	250	6	355	11	21
9400	0	1	250	5	352	3	17
9600	0	1	250	5	350	0	13
9800	0	1	249	1	347	0	10
10000	0	1	248	1	345	0	8

AMP SECONDS= 162,497196

VOLT SECONDS .089051

RUNID: CW1142 PROJECT: 02980

USER: 4126810219

ITEM	AMOUNT	COST
CPU TIME	00:00:14.630	\$0.43
FILE I/O REQUESTS	363	\$0.13
FILE I/O WORDS	302008	\$0.04
SWAP REQUESTS	1	\$0.00
SWAP WORDS	3584	\$0.00
MEMORY USAGE	0.148	\$0.10
CARDS IN	85	\$0.08
PAGES PRINTED	11	\$0.27
SHARED FILES	1	\$0.10
SOFTWARE SUPPORT	16	\$0.23
JOB CHARGE	1	\$0.05
TOTAL COST		\$1.43

THE ABOVE DOLLAR AMOUNTS ARE APPROXIMATE AND ARE BASED ON RATES FOR WHILE
USER BALANCE \$432.29

INITIATION TIME: 14:36:39 JAN 5, 1979
TERMINATION TIME: 14:41:54 JAN 5, 1979
PREVIOUS RUN TIME: 14:51:32 DEC 31, 1978

FIGURE 1

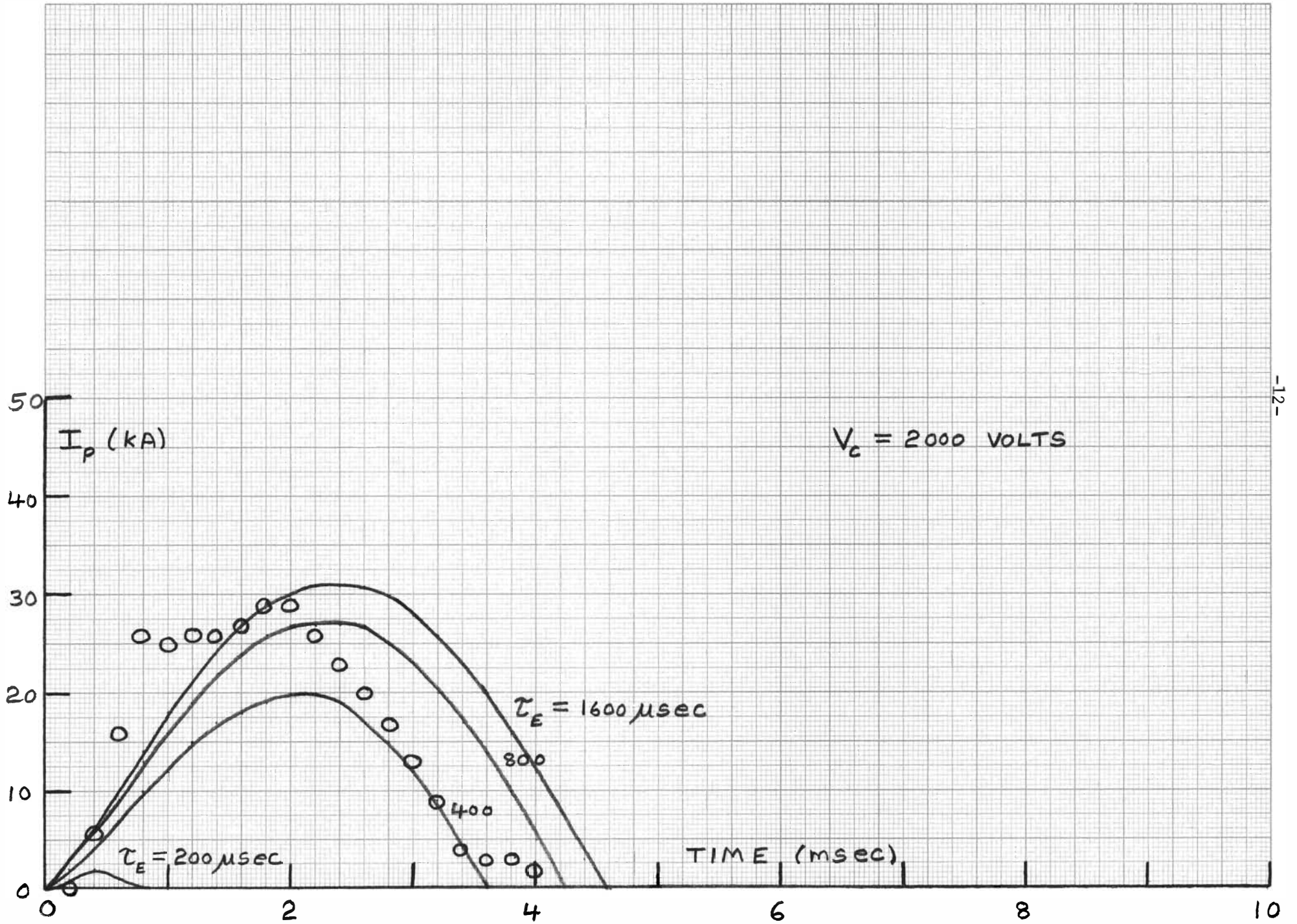


FIGURE 2

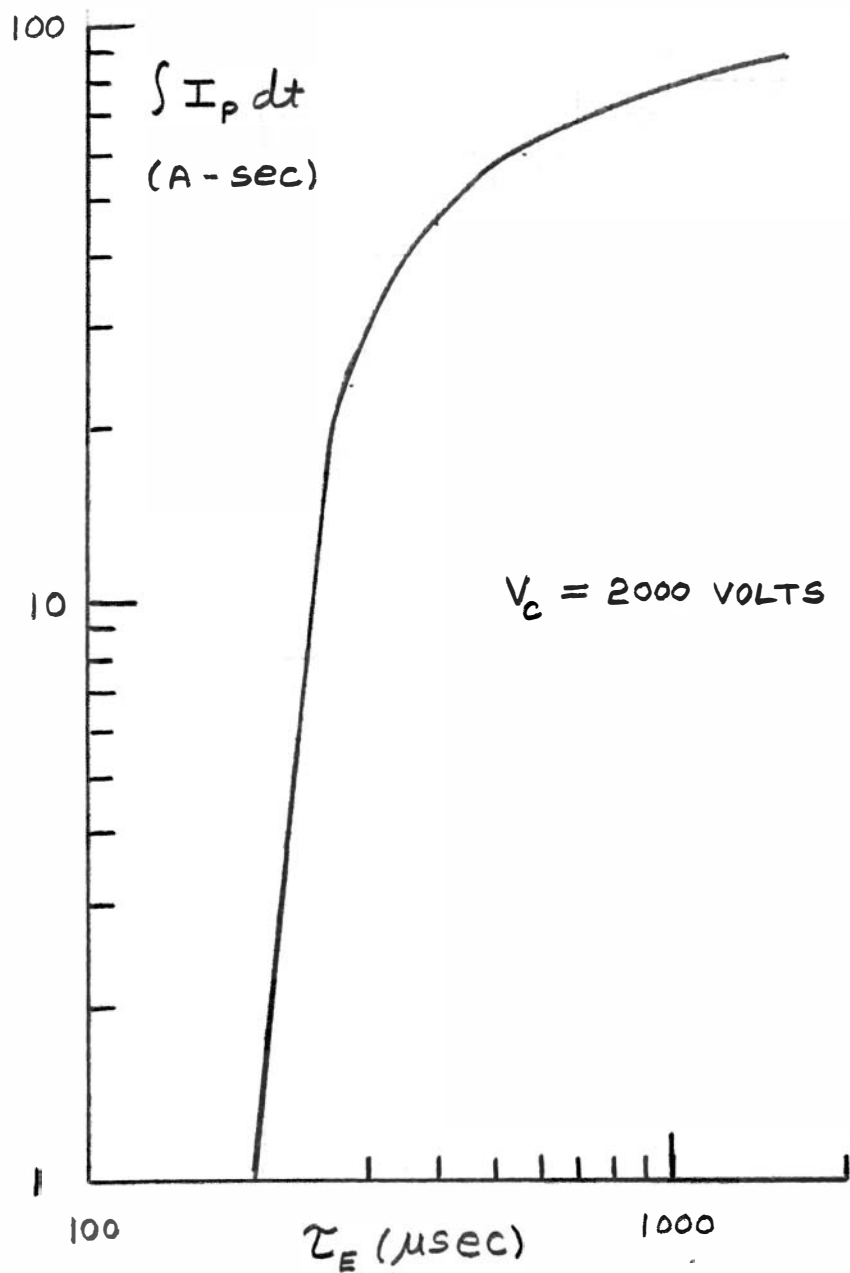


FIGURE 3

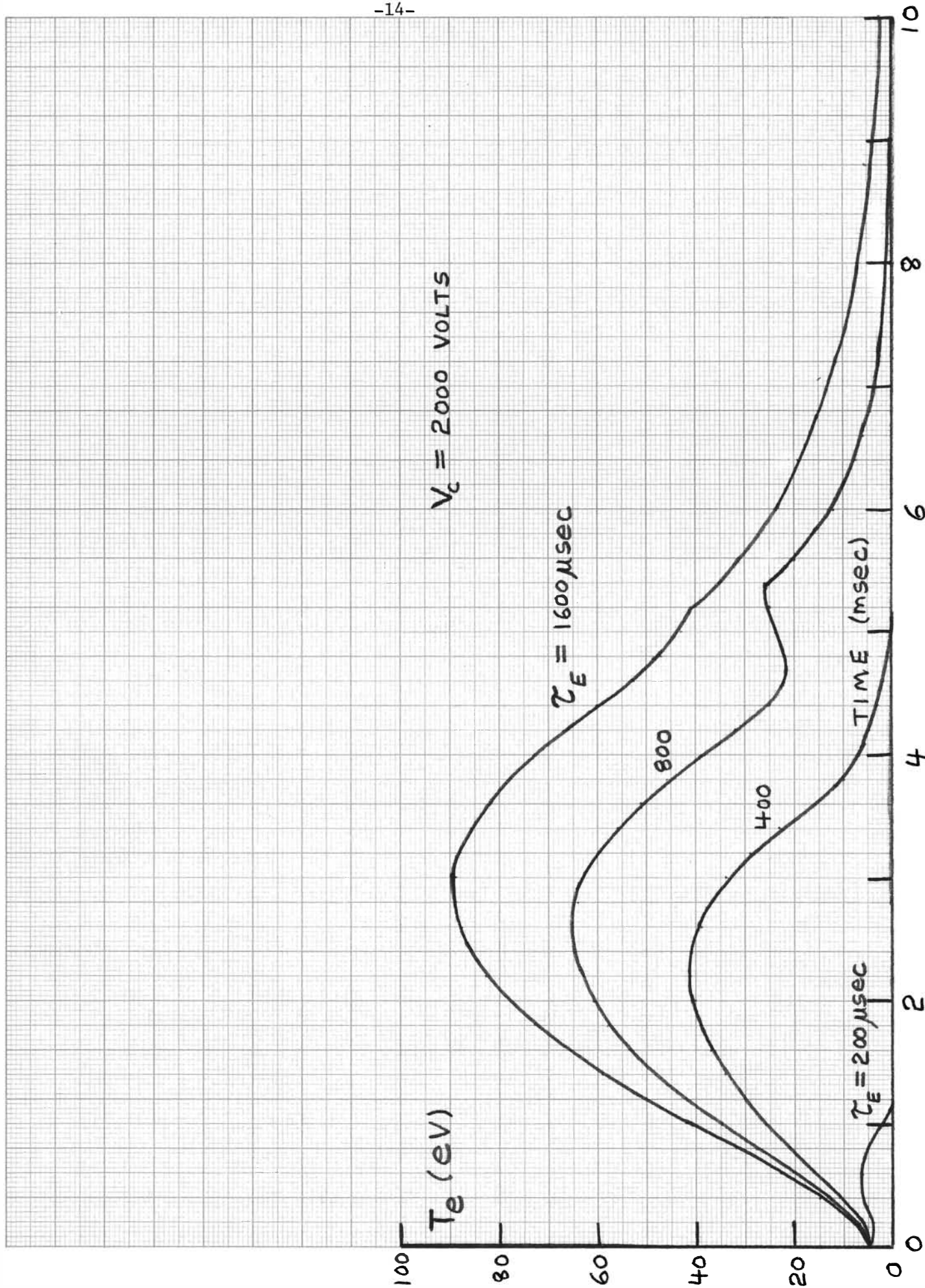
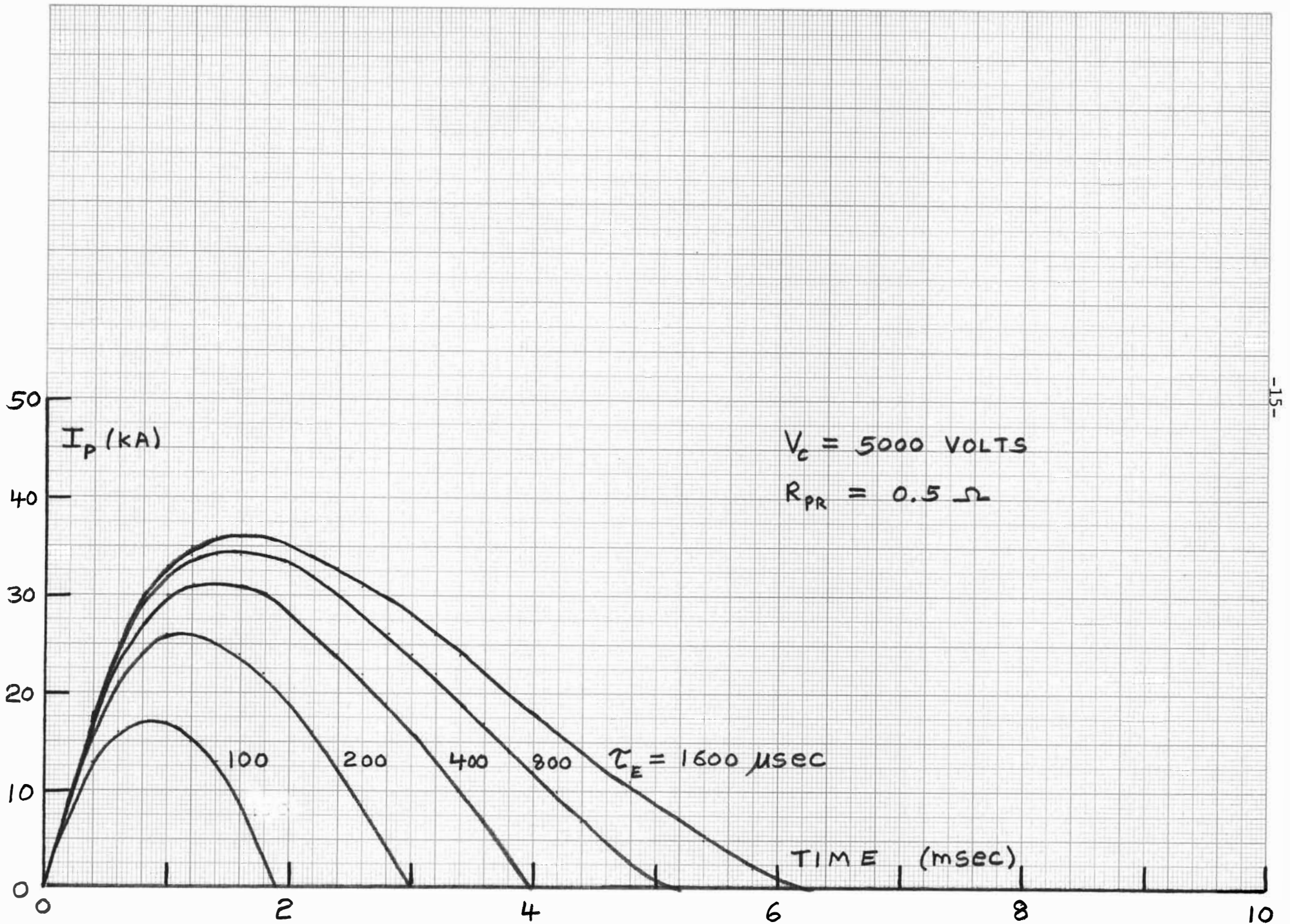


FIGURE 4



-15-

FIGURE 5

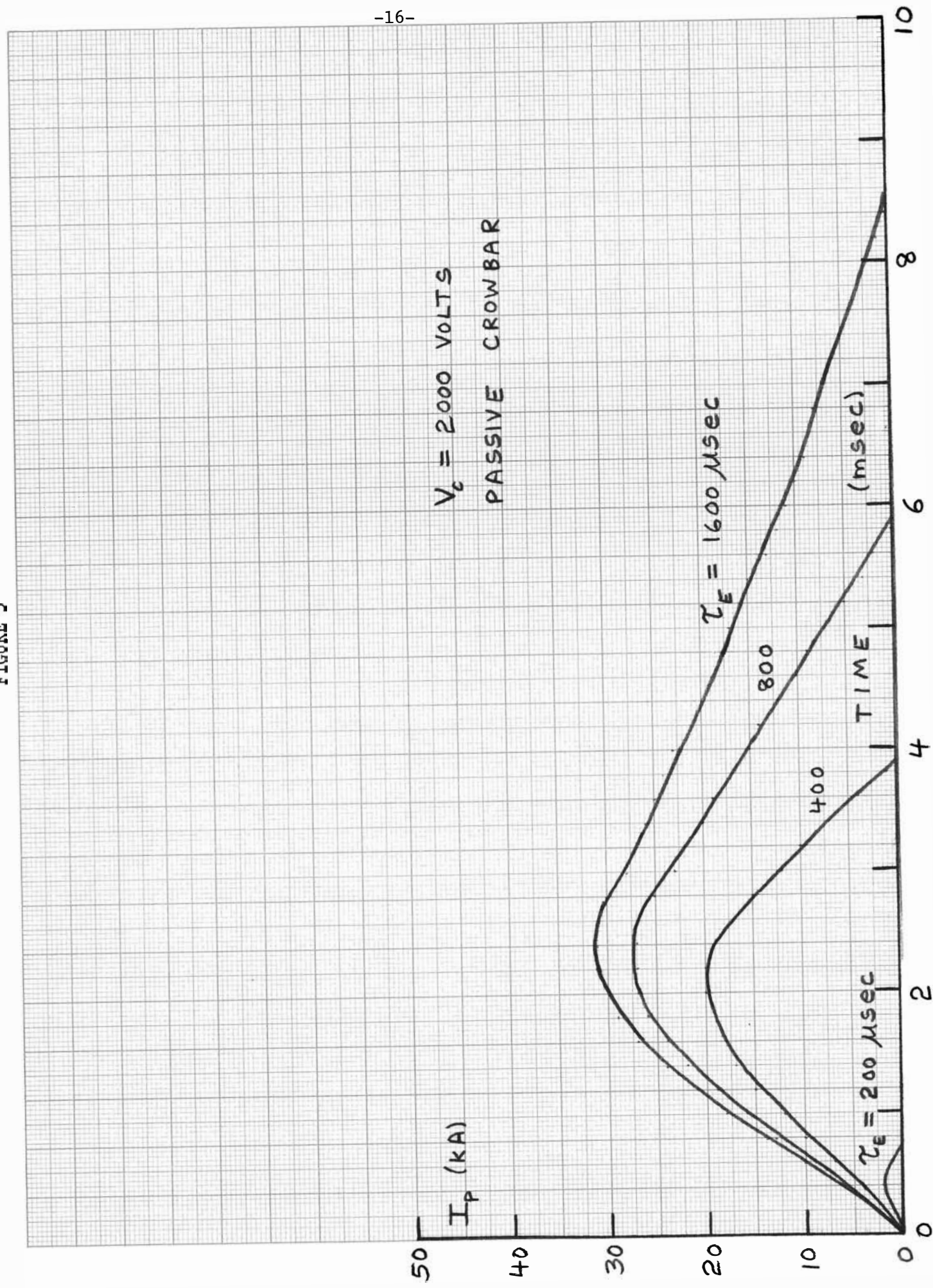


FIGURE 6

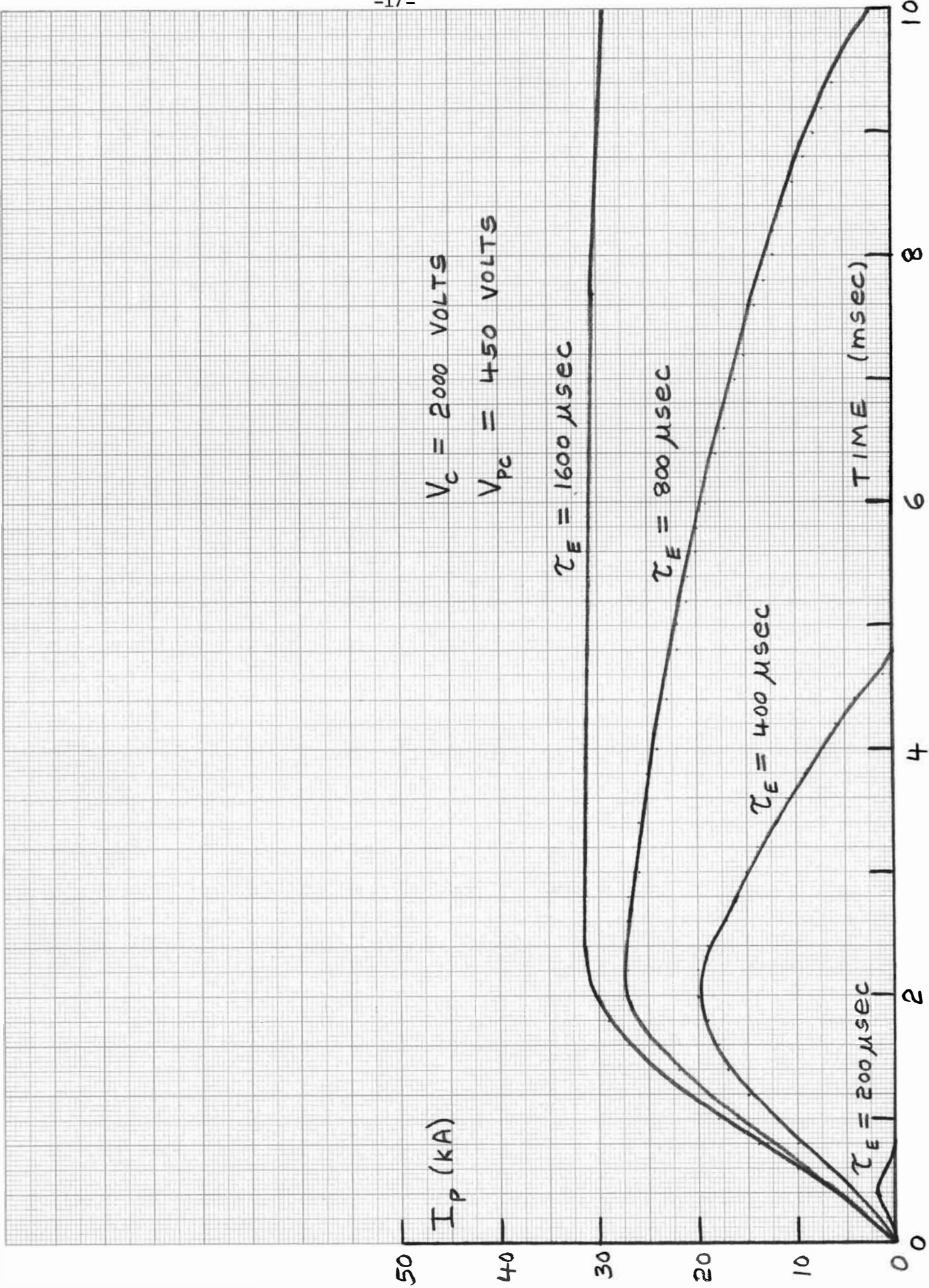


FIGURE 7

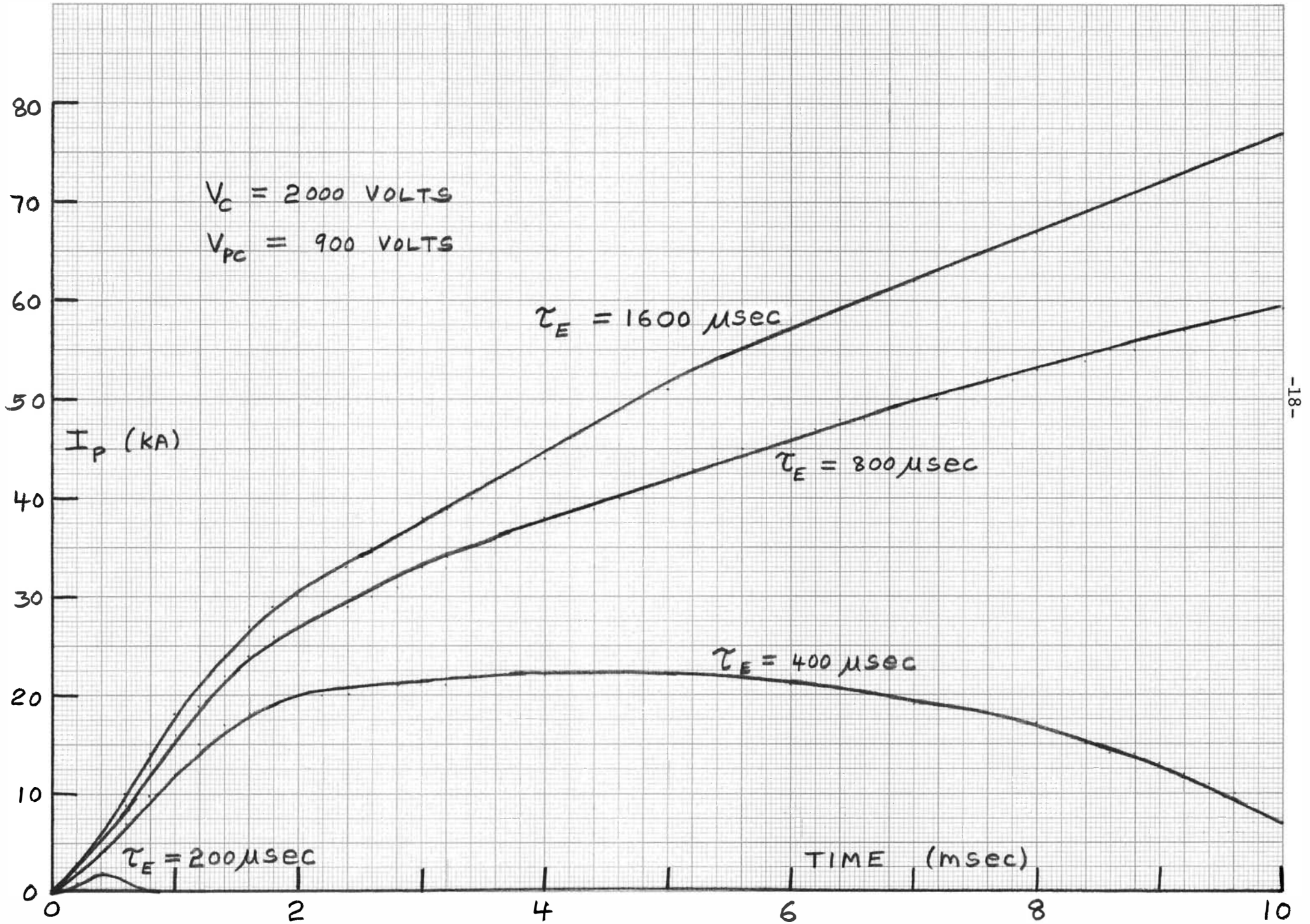


FIGURE 8

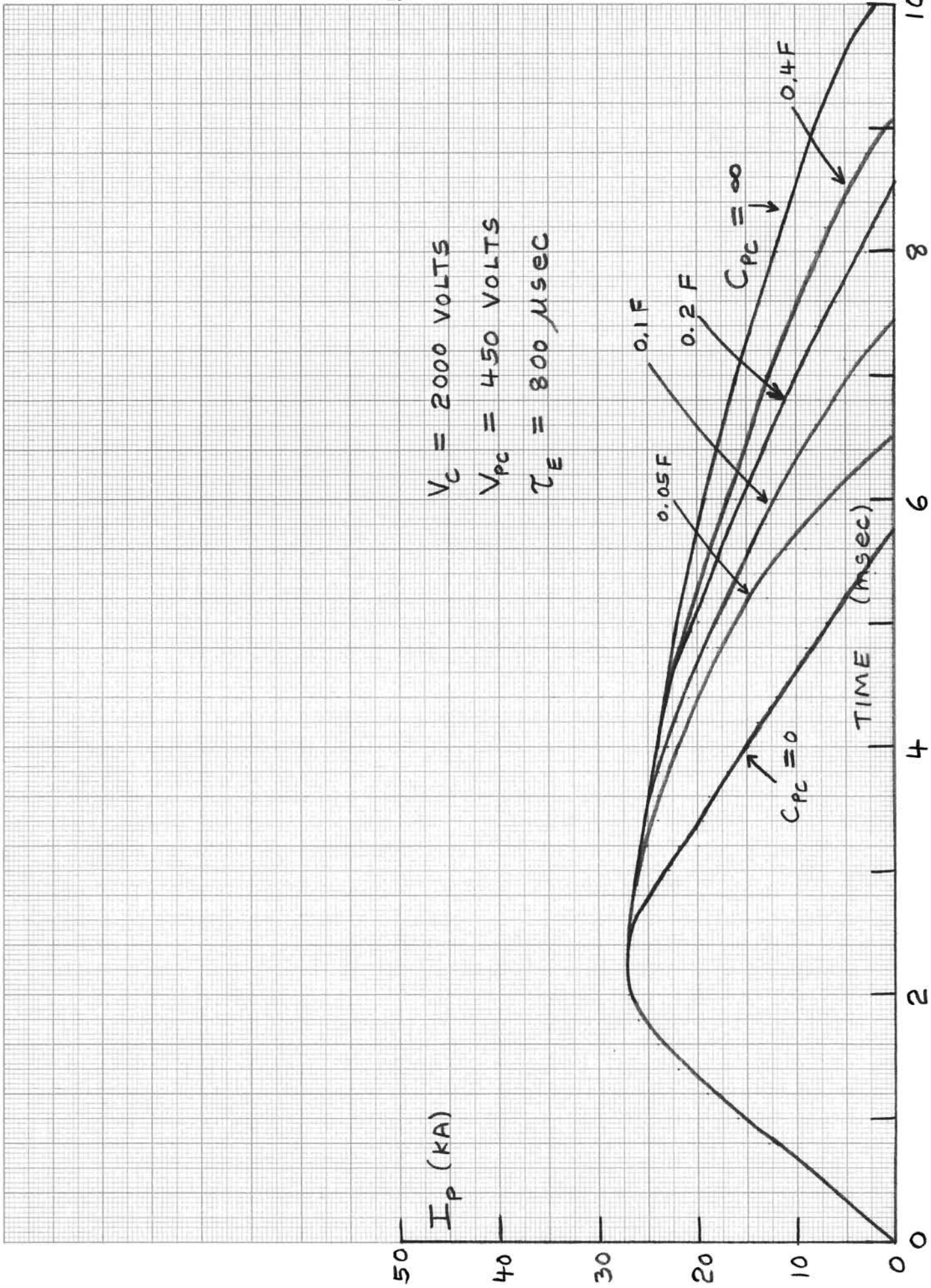


FIGURE 9

